# THE BIHARMONIC OPTIMAL SUPPORT PROBLEM 

ANTOINE LEMENANT AND MOHAMMAD REZA PAKZAD


#### Abstract

We establish a $\Gamma$-convergence result for $h \rightarrow 0$ of a thin nonlinearly elastic 3Dplate of thickness $h>0$ which is assumed to be glued to a support region in the 2D-plane $x_{3}=0$ over the $h$-2D-neighborhood of a given closed set $K$. In the regime of very small vertical forces we identify the $\Gamma$-limit as being the bi-harmonic energy, with Dirichlet condition on the gluing region $K$, following a general strategy by Friesecke, James, and Müller that we have to adapt in presence of the glued region. Then we introduce a shape optimization problem that we call "optimal support problem" and which aims to find the best glued plate. In this problem the bi-harmonic energy is optimized among all possible glued regions $K$ that we assume to be connected and for which we penalize the length. By relating the dual problem with Griffith almost-minimizers, we are able to prove that any minimizer is $C^{1, \alpha}$ regular outside a set of Hausdorff dimension strictly less then one.


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## 1. Introduction

Let $\Omega \subset \mathbb{R}^{2}$ a bounded Lipschitz domain, and $K \subset \bar{\Omega}$ be a closed set. For any $k \geq 1$ we consider the subspace of $W^{k, 2}(\Omega)$ defined by
$H_{0, K}^{k}(\Omega):=\left\{u \in W^{k, 2}(\Omega)\right.$ such that $\left.\partial^{\alpha} u\right|_{K}=0$ for all multiindices $\alpha$ such that $\left.0 \leq \alpha \leq k-1\right\}$, where by $\left.\partial^{\alpha} u\right|_{K}=0$ we mean that $\lim _{r \rightarrow 0}\left(f_{B_{r}\left(x_{0}\right) \cap \Omega}\left|\partial^{\alpha} u\right| d x\right)=0$ for $(k-|\alpha|, 2)$-q.e. $x_{0}$ on $K$.

It follows from the literature (see Lemma A. 1 in the appendix for a proof), that if $\Omega \subset \mathbb{R}^{N}$ is a Lipschitz domain then

$$
H_{0, K}^{k}(\Omega)=\left\{u \in W^{k, 2}(\Omega) \text { such that } \mathcal{E}(u) \in W_{0}^{k, 2}\left(\mathbb{R}^{N} \backslash K\right)\right\}
$$

where $\mathcal{E}$ is any extension operator $\mathcal{E}: W^{k, 2}(\Omega) \rightarrow W^{k, 2}\left(\mathbb{R}^{N}\right)$ and that $H_{0, K}^{k}(\Omega)$ stands for a closed subspace of $W^{k, 2}(\Omega)$.

Letting $f \in L^{2}(\Omega)$, this provides the existence of $u_{K}$ being the unique solution of

$$
\left\{\begin{array}{c}
\Delta^{2} u_{K}=f  \tag{1.1}\\
u \in H_{0, K}^{2}(\Omega)
\end{array}\right.
$$

In particular, $u_{K}$ is the unique critical point to the bi-harmonic energy

$$
\begin{equation*}
E(u):=\frac{1}{2} \int_{\Omega}\left|\nabla^{2} u\right|^{2} d x-\int_{\Omega} u f d x \tag{1.2}
\end{equation*}
$$

over $H_{0, K}^{2}(\Omega)$, and the associated compliance energy is

$$
\begin{equation*}
\int_{\Omega} u_{K} f d x=\int_{\Omega}\left|\nabla^{2} u_{K}\right|^{2} d x \tag{1.3}
\end{equation*}
$$

In this paper we intend to derive the functional (1.2) as the $\Gamma$-limit of a support-gluing problem for the 3d nonlinear thin elastic body subject to vertical body forces very small with respect to its thickness, which is the main result of the present paper.

To be more precise, we denote by $\Omega^{h}:=\Omega \times(0, h)$, which models a thin elastic plate which is assumed to be glued to a support in the plane $x_{3}=0$ by applying a surface glue to the support region

$$
K_{h}:=\left\{\left(x^{\prime}, 0\right) \in \Omega \times\{0\} ; \operatorname{dist}\left(x^{\prime}, K\right)<h\right\}
$$

Then we assume the plate to be subject to a vertical body force $\mathbf{f}^{h}: \Omega^{h} \rightarrow \mathbb{R}^{3}, \mathbf{f}^{h}:=\left(0,0, h^{\alpha} \tilde{f}\right)$, where $\alpha>2, \tilde{f}\left(x^{\prime}, x_{3}\right)=f\left(x^{\prime}\right)$ for $f \in L^{2}(\Omega)$.

In the most simplified framework (see Section 2.1 for a general case), the bulk elastic energy of a deformation $y: \Omega^{h} \rightarrow \mathbb{R}^{3}$ is given by

$$
E^{h}(y)=\frac{1}{h} \int_{\Omega^{h}} \frac{1}{2} \stackrel{2}{\operatorname{dist}}(\nabla y, S O(3)) d x
$$

The deformation of the plate $\Omega^{h}$ subject to body forces and gluing constraint is then variationaly modeled by minimizing the energy functional

$$
J^{h}(y):=E^{h}(y)-\frac{1}{h} \int_{\Omega^{h}} \mathbf{f}^{h} \cdot y d x
$$

among all deformations in the class

$$
\begin{equation*}
\mathcal{A}_{K}^{h}:=\left\{y \in W^{1,2}\left(\Omega^{h}, \mathbb{R}^{3}\right) ;\left.y\right|_{K_{h} \times\{0\}}=\operatorname{id}_{K_{h} \times\{0\}}\right\} \tag{1.4}
\end{equation*}
$$

The first goal of this paper is to establish the $\Gamma$-convergence of $J^{h}$ towards the bi-harmonic functional (1.2), in the regime of very small deformations in the sense that there exists $\beta>4$ such that

$$
\begin{equation*}
\limsup _{h \rightarrow 0} \frac{1}{h^{\beta}}\left(J^{h}\left(y^{h}\right)-\inf _{\mathcal{A}_{K}^{h}} J^{h}\right)=0 \tag{1.5}
\end{equation*}
$$

Here is a shorter version of one of the main statements of our paper, stated, moreover, for simplicity here in the introduction, for the most simplified functionals $E^{h}$ and $J^{h}$ as above. In order to state this result we introduce the average vertical displacement functions

$$
u_{h} \in W^{1,2}(\Omega, \mathbb{R}), \quad U_{h}:=\frac{1}{h} \int_{0}^{h}\left(y_{3}^{h}\left(\cdot, x_{3}\right)-h x_{3}\right) d x_{3}, \quad u_{h}:=\frac{1}{h^{\beta / 2-1}} U_{h}
$$

where $y_{3}^{h}$ is the vertical component of the deformation $y^{h}$.
Theorem 1.1 ( $\Gamma$-convergence). Assume that $\Omega \subset \mathbb{R}^{2}$ is a Lipschitz domain and $K \subset \bar{\Omega}$ is a closed set satisfying $\operatorname{Cap}_{1,2}(K)>0$. Assume furthermore that $\beta>4$. Then the functional $J^{h}$ $\Gamma$-converges to the energy in (1.2) in the following sense.
(1) for any sequence $h \rightarrow 0$ such that $y^{h} \in \mathcal{A}_{K}^{h}$ with energy bound $E^{h}\left(y^{h}\right) \precsim h^{\beta}$, there eone can find $u \in H_{0, K}^{2}(\Omega)$ and a subsequence such that $u_{h} \xrightarrow{\text { in } W^{1,2}} u \in W^{2,2}(\Omega)$, and

$$
\underset{h}{\liminf } \frac{1}{h^{\beta}} J^{h}\left(y^{h}\right) \geq \int_{\Omega} \frac{1}{24}\left|\nabla^{2} u\right|^{2}-u f d x
$$

(2) for every $u \in H_{0, K}^{2}(\Omega)$, there exists a sequence $y^{h} \in \mathcal{A}_{K}^{h}$ such that $u_{h} \xrightarrow{\text { in } W^{1,2}} u$ and

$$
\limsup _{h} \frac{1}{h^{\beta}} J^{h}\left(y^{h}\right) \leq \int_{\Omega} \frac{1}{24}\left|\nabla^{2} u\right|^{2}-u f d x
$$

The first part of Theorem 1.1 follows from Theorem 2.10 which gives a more general compactness result for sequences satisfying $E^{h}\left(y^{h}\right) \precsim h^{\beta}$. In Section 2.5 we explain how to obtain the liminf inequality in (1) (see Theorem 2.13). The second part is the purpose of Section 3. Notice that our main results are actually more general and work for a nonlinear energy of of the form $\int W(\nabla y) d x$ with $W$ satisfying some standard assumptions.

For the bulk of the argument, the proof or our result is inspired by the seminal papers by Friesecke, James, and Müller [14, 15], specially regarding the so called linearized von Kármán theory in [15, Theorem 2]. The novelty we need to take care of here is the "gluing" part of the problem on the support $K_{h}$ and its passage to the limit. Notice that we only assume $y=\mathrm{id}$ on the 2 D bottom part $K_{h} \subset \Omega \times\{0\}$ which differs from a standard Dirichlet or clamped boundary conditions. In particular, a key step to obtain compactness of a sequence with bounded energy in a thin plate is to be able to approximate it by piecewise constant maps with values in $S O(3)$. In our context, we have to adapt this approximation by insuring the map to be constant equal to identity on the support set $K_{h}$. As $K_{h}$ is a very thin $2 D$ set in $\mathbb{R}^{3}$, this part is non trivial and uses thin properties on Sobolev functions (see Corollary 2.9 for a statement).

Since our goal is to arrive to the bi-harmonic problem, we consider in this paper only the case of $\beta>4$, even if some other regimes could also be probably investigated. As a result, we can interpret the solution $u_{K}$ in (1.1) as modeling in the linear regime a vertical displacement of the 2D-plate $\Omega$, attached, or supported, onto the set $K$, and by (1.3) the energy associated to this displacement. Then, one can seek for the "best way" of attaching the plate, when the support
$K$ is penalized by its length. This leads to the following bi-harmonic optimal support problem:

$$
\begin{equation*}
\min _{K \subset \bar{\Omega}} \int_{\Omega}\left|\nabla^{2} u_{K}\right|^{2} d x+\mathcal{H}^{1}(K), \tag{1.6}
\end{equation*}
$$

where the minimum is taken over all compact and connected subsets $K$.
It is worth mentioning that replacing the bi-harmonic operator by the standard Laplace operator would lead to the so-called "optimal compliance problem" that was studied before in many papers (see $[5,9,10,21,22,6,8,7]$ ). A variant with a $p$-Laplacian has been also studied in $[6,8,7]$.

In other words, our paper deals with a bi-laplacian variant of the standard optimal compliance problem, which turns out to be relevant from the mechanical point of view. One of the motivation for this paper was indeed to provide a better justification of the so-called "optimal compliance problem". In the second part of the paper we prove existence and regularity results for a minimizer $K$ of problem (1.6), leading to a second result proved within this paper.

Theorem 1.2. For every bounded domain $\Omega \subset \mathbb{R}^{2}$ there exists a minimizer $K$ for the optimal support problem in (1.6). Moreover, if $K \cup \partial \Omega$ is connected, then any minimizer $K$ is an almost minimizer of the Griffith functional, and therefore is locally $C^{1, \alpha}$ regular inside $\Omega$, except for a singular set of points with Hausdorff dimension strictly less then 1.

To prove Theorem 1.2 we consider the dual formulation of the problem in (1.6), and prove that the minmiizers of the dual problem are almost minimizers of the so-called Griffith functional, already studied in the literature. By use the regularity results contained in [18] and [17], we obtain the conclusion.
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## 2. Justification of the biharmonic support model from 3d nonlinear elasticity

2.1. The 3d support-gluing model. As before in the introduction, let $\Omega \subset \mathbb{R}^{2}$ be a bounded Lipschitz domain, $K \subset \bar{\Omega}$ be a closed set and let $\Omega^{h}:=\Omega \times(0, h)$ a thin elastic plate. As before it is assumed that the thin plate is glued to a support in the plane $x_{3}=0$ by applying a surface glue to the support region

$$
K_{h}:=\left\{\left(x^{\prime}, 0\right) \in \Omega \times\{0\} ; \operatorname{dist}\left(x^{\prime}, K\right)<h\right\},
$$

and is subject to a vertical body force $\mathbf{f}^{h}: \Omega^{h} \rightarrow \mathbb{R}^{3}, \mathbf{f}^{h}:=\left(0,0, h^{\alpha} \tilde{f}\right)$, where $\alpha>2$, $\tilde{f}\left(x^{\prime}, x_{3}\right)=f\left(x^{\prime}\right)$ for $f \in L^{2}(\Omega)$. Here we have opted for a simplified set of assumptions on the body forces in order to focus on the new contribution.

In what follows

$$
S O(n):=\left\{R \in \mathbb{R}^{n \times n} ; R^{T} R=\mathrm{Id}, \operatorname{det} R>0\right\}
$$

is the special orthogonal group of 3d rotation matrices. The elastic density or potential $W$ : $\mathbb{R}^{3 \times 3} \rightarrow \mathbb{R}$ is assumed to satisfy the following natural conditions for all $F \in \mathbb{R}^{3 \times 3}$ :

- Normalization: $W(F) \geq 0, W(\mathrm{Id})=0$.
- Frame invariance: $\forall R \in S O(3), W(R F)=W(F)$.
- Non-degeneracy: $W(F) \geq c \operatorname{dist}^{2}(F, S O(3))$ for a constant $c>0$.

Moreover we will assume that $W$ is of class $C^{2}$ in a neighborhood of $S O(3)$. The bulk elastic energy of a deformation $y: \Omega^{h} \rightarrow \mathbb{R}^{3}$ is given by

$$
E^{h}(y)=\frac{1}{h} \int_{\Omega^{h}} W(\nabla y) d x .
$$

The deformation of the plate $\Omega^{h}$ subject to body forces and gluing constraint is then variationally modeled by minimizing the energy functional

$$
J^{h}(y):=E^{h}(y)-\frac{1}{h} \int_{\Omega^{h}} \mathbf{f}^{h} \cdot y d x,
$$

among all deformations in the class

$$
\begin{equation*}
\mathcal{A}_{K}^{h}:=\left\{y \in W^{1,2}\left(\Omega^{h}, \mathbb{R}^{3}\right) ;\left.y\right|_{K_{h} \times\{0\}}=\operatorname{id}_{K_{h} \times\{0\}}\right\} . \tag{2.1}
\end{equation*}
$$

Notice that we will still denote by $E^{h}$ and $J^{h}$ some functionals that was introduced before in the introduction in the particular case of $W(F)=\frac{1}{2} \operatorname{dist}^{2}(F, S O(3))$.

For $\beta>0$, we say the sequence $y^{h}$ is a $\beta$-minimizing sequence for $J^{h}$ whenever

$$
\limsup _{h \rightarrow 0} \frac{1}{h^{\beta}}\left(J^{h}\left(y^{h}\right)-\inf _{\mathcal{A}_{K}^{h}} J^{h}\right)=0 .
$$

2.2. The limiting 2D energy. Following [15] we introduce the linearized energy of second order

$$
Q_{3}(F):=\frac{\partial^{2} W}{\partial F^{2}}(I d)(F, F),
$$

and $Q_{2}: \mathbb{R}^{2 \times 2} \rightarrow \mathbb{R}$ defined by

$$
Q_{2}(G)=\min _{a \in \mathbb{R}^{3}} Q_{3}\left(G+a \otimes e_{3}+e_{3} \otimes a\right) .
$$

Under our assumptions, both forms are positive, semidefinite, convex and positive definite on symmetric matrices.

For the special case of isotropic elasticity, i.e. when $W(F R)=W(F)$ for all $R \in S O(3)$ and $F \in \mathbb{R}^{3 \times 3}$, it can be shown that

$$
Q_{3}(F)=2 \mu\left|\frac{F+F^{T}}{2}\right|^{2}+\lambda(\operatorname{Tr} F)^{2},
$$

and

$$
Q_{2}(G)=2 \mu\left|\frac{G+G^{T}}{2}\right|^{2}+\frac{2 \mu \lambda}{2 \mu+\lambda}(\operatorname{Tr} G)^{2}
$$

where $\mu>0$ and $\lambda \geq 0$. In particular if $W(F)=\frac{1}{2} \operatorname{dist}^{2}(F, S O(3))$, then $\mu=1 / 2$ and $\lambda=0$.

### 2.3. Friesecke-James-Müller rigidity estimate in presence of affine boundary condi-

 tions. In this section we present a corollary of the celebrated geometric rigidity estimate of Friesecke-James-Müller [14] which will be a necessary ingredient of the compactness argument. We first state this rigidity estimate:Theorem 2.1. [14, Theorem 3.1] Let $n \geq 2$ and $\Omega \subset \mathbb{R}^{n}$ be a bounded Lipschitz domain. Then there exists a constant $C=C(\Omega)$, such that for all mapping $u \in W^{1,2}\left(\Omega, \mathbb{R}^{n}\right)$, there is a rotation $R \in S O(n)$ with

$$
\begin{equation*}
\|\nabla y-R\|_{L^{2}(\Omega)}^{2} \leq C\|\operatorname{dist}(\nabla y, S O(n))\|_{L^{2}(\Omega)}^{2} \tag{2.2}
\end{equation*}
$$

We would like to identify the "average rotation" $R$ in (2.2) based on the boundary conditions applied to $y$ on portions of $\partial \Omega$. In particular, we expect that the rotation $R$ can be chosen to be uniformly the identity matrix Id for all mappings for which $\left.y\right|_{S}=\operatorname{id}_{S}$ for an open set $S \subset \partial \Omega$ :

Definition 2.2. We say that a matrix $F \in \mathbb{R}^{n \times n}$ is rank-one connected to $S O(n)$ whenever there exists $R \in S O(n)$ for which $\operatorname{rank}(F-R)=1$. Note that no element of $S O(n)$ itself does enjoy this property.

Corollary 2.3. Let $n \geq 2$ and $\Omega \subset \mathbb{R}^{n}$ be a bounded Lipschitz domain. For $F \in \mathbb{R}^{n \times n}$, and $S$ an open connected subset of $\partial \Omega$, we define

$$
\mathcal{A}_{S, F}:=\left\{y \in W^{1,2}\left(\Omega, \mathbb{R}^{n}\right) ;\left.y\right|_{S}(x)=F x\right\}
$$

Let $R_{F}$ be any closest element of $S O(n)$ to $F$.
(i) If $F$ is not rank-one connected to $S O(n)$, there exists $C=C(\Omega, S, F)$ such that for all $y \in \mathcal{A}_{S, F}$,

$$
\left\|\nabla y-R_{F}\right\|_{L^{2}(\Omega)}^{2} \leq C\|\operatorname{dist}(\nabla y, S O(n))\|_{L^{2}(\Omega)}^{2}
$$

(ii) Assume moreover that $S$ is inside no hyperplane of $\mathbb{R}^{n}$. Then, there exists $C=C(\Omega, S)$ such that for all $y \in \mathcal{A}_{S, F}$,

$$
\left\|\nabla y-R_{F}\right\|_{L^{2}(\Omega)}^{2} \leq C\|\operatorname{dist}(\nabla y, S O(n))\|_{L^{2}(\Omega)}^{2}
$$

Remark 2.4. The estimate can fail to be true if $\operatorname{rank}(F-R)=1$ and $S$ is a subset of a hyperplane. As a counter-example, take such $R, F$ in a manner that $R \neq R_{F}$, let $y=R x$, and choose $S$ such that $R x=F x$ on $S$.

Proof. The proof follows the same approach as in [12, Proposition 3.4]. We write the full details for the reader's convenience. In what follows the constant $C$ might differ from line to line but it will always depend on $\Omega$. Its dependance on $S$ or $F$ will be clarified on the way.

Consider an arbitrary point $p \in S \subset \mathbb{R}^{n}$ to be fixed later. Since the domain is Lipschitz, there is a hyperplane $P \ni p$ such that $\partial \Omega$ is locally the graph of a Lipschitz function over an open subset $U \subset P$ containing $p$. Let $V$ be the intersection of the projection of $S$ on $P$ with $U$. Fix a rotation $R_{0} \in S O(n)$ such that $\psi(x):=R_{0} x+p$ maps $\mathbb{R}^{n-1} \times\{0\}$ to $P$. We let $\widetilde{\Omega}:=\psi^{-1}(\Omega)$, $\tilde{S}:=\psi^{-1}(S) \subset \partial \widetilde{\Omega}$ and we fix $r>0$ for which $\psi\left(B_{r}(0) \times\{0\}\right) \subset V$, where $B_{r}(0)$ is the ball of radius $r$ in $\mathbb{R}^{n-1}$. Hence there exists a Lipschitz function $g: B_{r}(0) \rightarrow \mathbb{R}$ such that whose graph is an open subset of $\tilde{S}$. We define $\phi: B_{r}(0) \rightarrow \tilde{S}$ to be the graph parameterization of the portion of $\tilde{S}$ which is over $B_{r}(0)$ :

$$
\phi(z):=\left[\begin{array}{c}
z \\
g(z)
\end{array}\right]=\left[\begin{array}{c}
z \\
\tilde{g}(z)
\end{array}\right]+\left[\begin{array}{l}
0 \\
c
\end{array}\right],
$$

where

$$
\tilde{g}(z):=g(z)-f_{B_{r}(0)} g(z) d z, \quad c:=f_{B_{r}(0)} g(z) d z .
$$

In the context of part (ii) we can choose $p$ and $r$ such that $\tilde{g}$ is not a linear function over $B_{r}(0)$, for otherwise $S$, being connected, would become part of a hyperplane.

Now assume that $y \in \mathcal{A}_{S, F}$ is given and let

$$
E:=\|\operatorname{dist}(\nabla y, S O(n))\|_{L^{2}(\Omega)}^{2} .
$$

Applying Theorem 2.1 to $y$ on $\Omega$ we obtain that for a uniform constant $C>0$, there exists $R \in S O(n)$ for which

$$
\begin{equation*}
\|\nabla y-R\|_{L^{2}(\Omega)}^{2} \leq C E \tag{2.3}
\end{equation*}
$$

We claim that under assumptions of parts (i) or (ii), $y \in \mathcal{A}_{S, F}$ implies

$$
\begin{equation*}
|F-R|^{2} \leq C E \tag{2.4}
\end{equation*}
$$

with $C$ depending on the respective claimed variables. Let us observe that this is sufficient to conclude the proof of the corollary: Indeed, our claim yields in view of the definition of $R_{F}$ and the fact that $R \in S O(n)$ :

$$
\left|R_{F}-R\right| \leq\left|F-R_{F}\right|+|F-R| \leq 2|F-R| \leq C \sqrt{E}
$$

This, combined with (2.3) implies

$$
\left\|\nabla y-R_{F}\right\|_{L^{2}(\Omega)}^{2} \leq 2\left(\|\nabla y-R\|_{L^{2}(\Omega)}^{2}+\left|R_{F}-R\right|^{2} \mathcal{L}^{n}(\Omega)\right) \leq C E
$$

which is the desired estimate. Hence what remains is to prove (2.4).
Letting $y:=y \circ \psi$ on $\widetilde{\Omega}$, and in view of $\nabla \tilde{y}=(\nabla y \circ \psi) R_{0}$, we obtain from (2.3)

$$
\begin{equation*}
\left\|\nabla \tilde{y}-R R_{0}\right\|_{L^{2}(\tilde{\Omega})}^{2} \leq C E \tag{2.5}
\end{equation*}
$$

We let $b:=f_{\Omega}\left(\tilde{y}(x)-R R_{0} x\right) d x$. Applying the Poincaré inequality on $\widetilde{\Omega}$, and the trace embedding of $W^{1,2}(\widetilde{\Omega})$ into $L^{2}(\partial \widetilde{\Omega})$, we obtain

$$
\left\|\tilde{y}-\left(R R_{0} x+b\right)\right\|_{L^{2}(\tilde{S})}^{2} \leq C\left\|\tilde{y}-\left(R R_{0} x+b\right)\right\|_{W^{1,2}(\tilde{\Omega})}^{2} \leq C E .
$$

But $\left.\tilde{y}\right|_{\tilde{S}}=F\left(R_{0} x+p\right)=F R_{0} x+F p$. This yields

$$
\begin{equation*}
\left\|(F-R) R_{0} x+(F p-b)\right\|_{L^{2}(\tilde{S})}^{2} \leq C E \tag{2.6}
\end{equation*}
$$

In what follows, and for any $F \in \mathbb{R}^{n \times n}$, we will denote by $\widehat{F} \in \mathbb{R}^{n \times(n-1)}$ and $F^{(n)} \in \mathbb{R}^{n}$, respectively the matrix made by the first $n-1$ columns of $F$, and the last column of $F$. (2.6) gives

$$
\begin{equation*}
e:=\left\|(F-R) R_{0} x+(F p-b)\right\|_{L^{2}\left(\phi\left(B_{r}(0)\right)\right.}^{2} \leq\left\|(F-R) R_{0} x+(F p-b)\right\|_{L^{2}(\tilde{S})}^{2} \leq C E . \tag{2.7}
\end{equation*}
$$

We estimate from below the left hand side through change of variable:

$$
\begin{aligned}
e & =\int_{\phi\left(B_{r}(0)\right)}\left|(F-R) R_{0} x+(F p-b)\right|^{2} d \mathcal{H}^{n-1} \\
& =\int_{B_{r}(0)}\left|(F-R) R_{0} \phi(z)+(F p-b)\right|^{2} \sqrt{1+|\nabla \phi|^{2}(z)} d z \\
& \geq \int_{B_{r}(0)}\left|(F-R) R_{0} \phi(z)+(F p-b)\right|^{2} d z \\
& =\int_{B_{r}(0)}\left|(F \widehat{-R}) R_{0} z+g(z)\left((F-R) R_{0}\right)^{(n)}+(F p-b)\right|^{2} d z \\
& =\int_{B_{r}(0)} \mid\left(\widehat{F-R)} R_{0} z+\tilde{g}(z)\left((F-R) R_{0}\right)^{(n)}+c\left((F-R) R_{0}\right)^{(n)}+\left.(F p-b)\right|^{2} d z\right.
\end{aligned}
$$

We let

$$
A:=\left(\widehat{F-R)} R_{0}, b^{\prime}:=\left((F-R) R_{0}\right)^{(n)}, \tilde{b}:=c b^{\prime}+(F p-b) .\right.
$$

Expanding the right hand side yields

$$
\begin{aligned}
e & =\int_{B_{r}(0)}|A z|^{2} d z+\left|b^{\prime}\right|^{2} \int_{B_{r}(0)}|\tilde{g}(z)|^{2} d z+|\tilde{b}|^{2}\left|B_{r}(0)\right| \\
& +2 \int_{B_{r}(0)} \tilde{g}(z)\left\langle A z, b^{\prime}\right\rangle d z+2 \int_{B_{r}(0)}\langle A z, \tilde{b}\rangle+\tilde{g}(z)\left\langle b^{\prime}, \tilde{b}\right\rangle d z \\
& \geq \int_{B_{r}(0)}|A z|^{2} d z+\left|b^{\prime}\right|^{2} \int_{B_{r}(0)}|\tilde{g}(z)|^{2} d z+2\left\langle\int_{B_{r}(0)} \tilde{g}(z) A z d z, b^{\prime}\right\rangle \\
& +2\left\langle\int_{B_{r}(0)} z d z, A^{T} \tilde{b}\right\rangle+2 \int_{B_{r}(0)} \tilde{g}(z) d z\left\langle b^{\prime}, \tilde{b}\right\rangle .
\end{aligned}
$$

To proceed we use the facts that $\int_{B_{r}(0)} z d z=0$ and $\int_{B_{r}(0)} \tilde{g}(z) d z=0$ to obtain

$$
\begin{equation*}
e \geq \int_{B_{r}(0)}|A z|^{2} d z+\left|b^{\prime}\right|^{2} \int_{B_{r}(0)}|\tilde{g}(z)|^{2} d z+2\left\langle\int_{B_{r}(0)} A z d z, \tilde{g}(z) b^{\prime}\right\rangle \tag{2.8}
\end{equation*}
$$

Now, we observe that there exists $0 \leq \rho<1$ such that

$$
\left|\left\langle\int_{B_{r}(0)} A z d z, \tilde{g}(z) b^{\prime}\right\rangle\right| \leq \rho\left|b^{\prime}\right|\left(\int_{B_{r}(0)}|\tilde{g}(z)|^{2} d z\right)^{\frac{1}{2}}\left(\int_{B_{r}(0)}|A z|^{2} d z\right)^{\frac{1}{2}} .
$$

If either $b^{\prime}=0$ or $A=0$ the claim is obvious. Otherwise if no such $\rho$ exists, by the CauchySchwartz inequality the two vector valued functions $\tilde{g}(z) b^{\prime}$ and $A z$ must be co-linear over $B_{r}(0)$ and for a $\lambda \neq 0$ we have

$$
\tilde{g}(z) b^{\prime}=\lambda A z
$$

Since $b^{\prime} \neq 0$, this implies that $\tilde{g}$ is linear and hence $B(x, r)$ must be a portion of a hyperplane, in contradiction to our original choice for part (ii). So in this case as $\tilde{g}$ is not affine, $\rho$ depends only on the local nonlinearity of $g$, irrespective of what $A$ is, and thus can be globally measured by how far $S$ is from being a hyperplane. If otherwise $g$ is affine, under the assumptions of part (i), it is $F$ which determines a value for $\rho$ independent of $\tilde{g}$. Indeed, we observe that $(F-R) R_{0} e_{j}=A e_{j}=\lambda^{-1} \tilde{g}\left(e_{j}\right) b^{\prime}$ for $j=1, \cdots, n-1$, and $(F-R) R_{0} e_{n}=b^{\prime}$ by definition of $b^{\prime}$, which implies that $F-R$ is of rank 1 , contrary to our assumption.

Following the above observation, (2.7) and (2.8) imply

$$
\begin{aligned}
& \int_{B_{r}(0)}|A z|^{2} d z+\left|b^{\prime}\right|^{2} \int_{B_{r}(0)}|\tilde{g}(z)|^{2} d z \\
& \leq \frac{1}{1-\rho}\left(\int_{B_{r}(0)}|A z|^{2} d z+\left|b^{\prime}\right|^{2} \int_{B_{r}(0)}|\tilde{g}(z)|^{2} d z-2 \rho\left|b^{\prime}\right|\left(\int_{B_{r}(0)}|\tilde{g}(z)|^{2} d z\right)^{\frac{1}{2}}\left(\int_{B_{r}(0)}|A z|^{2} d z\right)^{\frac{1}{2}}\right) \\
& \leq \frac{1}{1-\rho} e \leq \frac{C}{1-\rho} E .
\end{aligned}
$$

Defining $G:=A^{T} A$, we first observe that

$$
|A|^{2}=\operatorname{Tr}(G)
$$

On the other hand, since $G$ is symmetric nonnegative, $G=R_{1}^{T} D R_{1}$ with respectively diagonal matrix $D=\operatorname{diag}\left\{K_{1}, \cdots, K_{n-1}\right\}$ and orthogonal matrix $R_{1} \in S O(n-1)$, and $\operatorname{Tr}(G)=\operatorname{Tr}(D)$.

Hence, by symmetry,

$$
\begin{aligned}
\operatorname{Tr}(G) \int_{B_{r}(0)} z_{1}^{2} d z & =\sum_{j=1}^{n-1} \int_{B_{r}(0)} K_{j} z_{j}^{2} d z=\int_{B_{r}(0)}\langle z, D z\rangle d z=\int_{B_{r}(0)}\left\langle R_{1} z, D R_{1} z\right\rangle d z \\
& =\int_{B_{r}(0)}\langle z, G z\rangle d z=\int_{B_{r}(0)}|A z|^{2} d z
\end{aligned}
$$

Therefore we obtain

$$
|A|^{2}+\left|b^{\prime}\right|^{2} \int_{B_{r}(0)}|\tilde{g}(z)|^{2} d z \leq \frac{C}{(1-\rho) r^{n+2}} E .
$$

To see part (i), observe that if $S$ is no entirely flat, we can choose $p$ in such a manner that

$$
\int_{B_{r}(0)}|\tilde{g}(z)|^{2} d z>0 .
$$

Therefore we conclude with (2.4) and finish the proof:

$$
|F-R|^{2}=\left|(F-R) R_{0}\right|^{2}=|A|^{2}+\left|b^{\prime}\right|^{2} \leq C(\Omega, S) E .
$$

If, $S$ is inside a hyperplane, and $\tilde{g}$ is still affine, but non-zero, the same conclusion holds, but this time with $C$ depending also on $F$ through $\rho$. Finally, if $\tilde{g} \equiv 0$, the best we obtain is

$$
\left|\widehat{(F-R)} R_{0}\right|^{2} \leq \frac{C(\Omega, F)}{r^{n+2}} E .
$$

But in this case $F$ is assumed not to be rank one connected to $S O(n)$, which means that for all $\tilde{R} \in S O(n), m:=\operatorname{rank}\left((F-\tilde{R}) R_{0}\right)$ is either 0 or at least 2 . If $m=0$ for some $\tilde{R}$, then $F=\tilde{R} \in S O(n)$, in which case we note that

$$
|F-R|^{2}=\left|(F-R) R_{0}\right|^{2}=\left|(\tilde{R}-R) R_{0}\right|^{2} \leq n^{2} \left\lvert\,\left(\left.\tilde{R-R)} R_{0}\right|^{2}=n^{2} \left\lvert\,\left(\left.\widehat{F-R)} R_{0}\right|^{2} \leq \frac{C(\Omega, F)}{r^{n+2}} E,\right.\right.\right.\right.
$$

establishing (2.4), where we used the fact that for any rotation,

$$
R=\left[v_{1}, \cdots, v_{n}\right] \in S O(n),
$$

the last column is the exterior product

$$
v_{n}=v_{1} \wedge v_{2} \wedge \cdots \wedge v_{n-1}
$$

of the first $n-1$ columns, which yields for any two rotations $R_{1}, R_{2} \in S O(n)$ :

$$
\left|R_{1}^{(n)}-R_{2}^{(n)}\right| \leq(n-1)\left|\widehat{R_{1}}-\widehat{R_{2}}\right|,
$$

finally implying

$$
\left|R_{1}-R_{2}\right|^{2}=\left|\left(R_{1}-R_{2}\right)^{(n)}\right|^{2}+\left|\widehat{R_{1}-R_{2}}\right|^{2} \leq n^{2}\left|\widehat{R_{1}-R_{2}}\right|^{2} .
$$

If on the other hand, $m \geq 2$, it is straightforward to see that

$$
C_{F}:=\inf _{\tilde{R} \in S O(n)}\left|\widehat{(F-\tilde{R})} R_{0}\right|^{2}>0
$$

We therefore obtain

$$
\left.\left|F-R_{F}\right|^{2}=\frac{1}{C_{F}} \operatorname{dist}^{2}(F, S O(n)) C_{F} \leq \frac{C_{F}^{\prime}}{C_{F}} \right\rvert\,\left(\left.\widehat{F-R)} R_{0}\right|^{2} \leq \frac{C(\Omega, F)}{r^{n+2}} E,\right.
$$

and

$$
\left|R_{F}-R\right|^{2} \leq 4 n \leq \frac{C}{C_{F}} \left\lvert\,\left(\left.\widehat{F-R)} R_{0}\right|^{2} \leq \frac{C(\Omega, F)}{r^{n+2}} E,\right.\right.
$$

which once again completes the proof.

Corollary 2.5. Let $Q=(0,1)^{3}$ be the unit cube. For $0<r_{0}<1$, let

$$
S\left(x^{\prime}, r_{0}\right):=B\left(x^{\prime}, r_{0}\right) \cap(0,1)^{2},
$$

and

$$
\mathcal{A}_{r_{0}}:=\left\{y \in W^{1,2}\left(Q, \mathbb{R}^{3}\right) ;\left.\exists x^{\prime} \in(0,1)^{2} y\right|_{S\left(x^{\prime}, r_{0}\right)}(x)=x\right\} .
$$

Then there exists $C=C\left(r_{0}\right)$ such that for all $y \in \mathcal{A}_{r_{0}}$,

$$
\|\nabla y-\operatorname{Id}\|_{L^{2}(Q)}^{2} \leq C\|\operatorname{dist}(\nabla y, S O(n))\|_{L^{2}(Q)}^{2}
$$

### 2.4. Compactness for bounded sequences.

Lemma 2.6. Let $\Omega, \Omega^{h}$ and $K$ be as above. Then there exist constants $h_{0}>0, C>0,0<c<1$, depending only on $\Omega$ and $K$, such that for all $h<h_{0}$ and $y \in \mathcal{A}_{K}^{h}$ there exists a matrix valued mapping $\widetilde{F} \in W^{1,2}\left(\Omega, \mathbb{R}^{3 \times 3}\right)$, extended trivially to $\Omega^{h}$, for which the following estimates hold true:

$$
\begin{equation*}
\frac{1}{h}\|\nabla y-\widetilde{F}\|_{L^{2}\left(\Omega^{h}\right)}^{2} \leq C E^{h}(y), \quad\|\nabla \widetilde{F}\|_{L^{2}(\Omega)}^{2} \leq \frac{C}{h^{2}} E^{h}(y) \tag{2.9}
\end{equation*}
$$

Moreover, if

$$
\tilde{d}\left(x^{\prime}\right):= \begin{cases}\left|\widetilde{F}\left(x^{\prime}\right)-\mathrm{Id}\right| & \text { if } x^{\prime} \in K_{c h} \\ \operatorname{dist}\left(\widetilde{F}\left(x^{\prime}\right), S O(3)\right) & \text { otherwise }\end{cases}
$$

then

$$
\begin{equation*}
\|\tilde{d}\|_{L^{2}(\Omega)}^{2} \leq C E^{h}(y), \quad\|\tilde{d}\|_{L^{\infty}(\Omega)}^{2} \leq \frac{C}{h^{2}} E^{h}(y) . \tag{2.10}
\end{equation*}
$$

Proof. The proof closely follows [15, Theorem 6]. We will need only to make minor but careful adjustments using Corollary 2.5.

We cover $\Omega$ with open sets $\left\{U_{j}\right\}_{j=0}^{N}$ such that $\bar{U}_{0} \subset \Omega$ and $\partial \Omega$, and for $j=1, \cdots, N, U_{j} \cap \Omega$ is such that for an open interval $I_{j} \subset \mathbb{R}$ and a Lipschitz function $g_{j}: I_{j} \rightarrow \mathbb{R}$ we have

$$
U_{j} \cap \Omega=\left\{x \in U_{j}, x_{1} \in I_{j}, x_{2}>g_{j}\left(x_{1}\right)\right\} \quad \text { and } \quad U_{j} \cap \partial \Omega=\left\{x \in U_{j}, x_{1} \in I_{j}, x_{2}=g_{j}\left(x_{1}\right)\right\}
$$

for a suitable orthonormal coordinate system adapted to $U_{j} \cap \Omega$. We also consider the flattening bi-Lipschitz change of variable $\Phi_{j}: U_{j} \cap \bar{\Omega} \rightarrow \overline{\mathbb{R}_{+}^{2}}$ defined by $\Phi\left(x_{1}, x_{2}\right)=\left(x_{1}, x_{2}-g_{j}\left(x_{1}\right)\right)$. Let also $\theta_{j} \in C_{c}^{\infty}\left(\mathbb{R}^{2}\right)$ be a partition of unity subject to the family $U_{j}$, i.e.for all $j \in\{0,1, \cdots, N\}$,

$$
E_{j}:=\operatorname{supp} \theta_{j} \subset U_{j} \quad \text { and } \quad \sum_{j=0}^{N} \theta_{j}=1 .
$$

We let for all $x^{\prime} \in \Omega$

$$
\bar{F}\left(x^{\prime}\right):=\frac{1}{h} \int_{0}^{h} \nabla y\left(x^{\prime}, x_{3}\right) d x_{3} .
$$

Step 1. Interior local estimates: We will first construct a matrix field $\widetilde{F}_{0}$ on $E_{0} \subset U_{0}$ with useful local estimates.

We choose $h_{0}$ small enough such that $\operatorname{dist}\left(E_{0}, \partial U_{0}\right)>3 h_{0}$. For $S_{0,1}=(0,1)^{2}$ being the open unit square in $\mathbb{R}^{2}$, we consider a standard mollifier $\phi \in C_{c}^{\infty}(Q(0,1))$, with $\phi \geq 0$ and $\int_{S_{0,1}} \phi=1$, and we set $\phi_{h}(x):=h^{-2} \phi(x / h)$. and we define for each $x^{\prime} \in E_{0}$,

$$
\widetilde{F}_{0}\left(x^{\prime}\right):=\phi_{h} * \bar{F}\left(x^{\prime}\right)=\frac{1}{h^{3}} \int_{S_{x^{\prime}, h} \times(0, h)} \phi\left(\frac{x^{\prime}-z^{\prime}}{h}\right) \nabla y(z) d z^{\prime} d z_{3},
$$

where $z=\left(z^{\prime}, z_{3}\right) \in \mathbb{R}^{3}$, and $S_{x^{\prime}, h}=\left(x^{\prime}, x^{\prime}+h\right)^{2}$ is the square of edge size $h^{\prime}$ with its lower left corner at $x^{\prime}$.

We now consider the cube $Q=\mathcal{Q}_{0}\left(x^{\prime}, h\right)=S_{x^{\prime}, h} \times(0, h)$ and apply Theorem 2.1 to $\left.y\right|_{Q}$. Therefore there exists $R_{x^{\prime}} \in S O(n)$ such that

$$
\begin{equation*}
\left\|\nabla y-R_{x^{\prime}}\right\|_{L^{2}\left(\mathcal{Q}_{0}\left(x^{\prime}, h\right)\right)}^{2} \leq C \int_{\mathcal{Q}_{0}\left(x^{\prime}, h\right)} \operatorname{dist}^{2}(\nabla y, S O(3)) \tag{2.11}
\end{equation*}
$$

Moreover, fixing a constant $0<c_{0}<1$, in view of Corollary $2.5, R_{x^{\prime}}$ can be chosen to be equal to the identity matrix Id with a uniform constant $C>0$ in (2.11) if $K_{c_{0} h} \cap \overline{S_{x^{\prime}, h}} \neq \emptyset$. Indeed, in that case, there exists $\tilde{x}^{\prime} \in K_{c_{0} h} \cap S_{x^{\prime}, h}$, such that setting $r_{0}=1-c_{0}$ we have $y=\mathrm{id}$ on

$$
B\left(\tilde{x}^{\prime}, r_{0} h\right) \cap S_{x^{\prime}, h} \subset K_{h} \cap S_{x^{\prime}, h} \subset \partial Q,
$$

and Corollary 2.5 applies after a proper translation and rescaling. Also note that in this case $C$ is independent of $h$ as the boundary conditions and the cube estimates in both Theorem 2.1 and Corollary 2.5 are invariant under dilations and translations. Now, since

$$
\frac{1}{h^{3}} \int_{\mathcal{Q}_{0}\left(x^{\prime}, h\right)} \phi\left(\frac{x^{\prime}-z^{\prime}}{h}\right) d z=1
$$

letting $d \mu=h^{-3} \phi\left(\left(x^{\prime}-z^{\prime}\right) / h\right) d z$ and applying Jensen's inequality yields

$$
\left|\widetilde{F}_{0}\left(x^{\prime}\right)-R_{x^{\prime}}\right|^{2}=\left|\int_{\mathcal{Q}_{0}\left(x^{\prime}, h\right)}\left(\nabla y-R_{x^{\prime}}\right) d \mu\right|^{2} \leq \int_{\mathcal{Q}_{0}\left(x^{\prime}, h\right)}\left|\nabla y-R_{x^{\prime}}\right|^{2} d \mu,
$$

which implies, using a uniform bound on $\phi$ and the rigidity estimate,

$$
\begin{equation*}
\left|\widetilde{F}_{0}\left(x^{\prime}\right)-R_{x^{\prime}}\right|^{2} \leq \frac{C}{h^{3}} \int_{\mathcal{Q}_{0}\left(x^{\prime}, h\right)} \operatorname{dist}^{2}(\nabla y, S O(3)) . \tag{2.12}
\end{equation*}
$$

To obtain a bound on $\nabla \widetilde{F}_{0}$ we proceed in a similar manner. For all $\tilde{x}^{\prime} \in S_{x^{\prime}, h}$ we note that

$$
\int_{\mathcal{Q}_{0}\left(x^{\prime}, h\right)} \nabla \phi\left(\frac{\tilde{x}^{\prime}-z^{\prime}}{h}\right) d z=0,
$$

and we obtain this time using Cauchy-Schwarz inequality and a uniform bound on $\nabla \phi$

$$
\begin{aligned}
\left|\nabla \widetilde{F}_{0}\left(\tilde{x}^{\prime}\right)\right|^{2} & =\left|h^{-4} \int_{\mathcal{Q}_{0}\left(x^{\prime}, h\right)}\left(\nabla y-R_{\tilde{x}^{\prime}, h}\right) \nabla \phi\left(\frac{\tilde{x}^{\prime}-z^{\prime}}{h}\right) d z\right|^{2} \\
& \leq\left(\int_{\mathcal{Q}_{0}\left(x^{\prime}, h\right)}\left|\nabla y-R_{\tilde{x}^{\prime}, h}\right|^{2} d z\right) \int_{\mathcal{Q}_{0}\left(x^{\prime}, h\right)}\left|h^{-4} \nabla \phi\left(\frac{\tilde{x}^{\prime}-z^{\prime}}{h}\right)\right|^{2} d z
\end{aligned}
$$

yielding

$$
\begin{align*}
\forall \tilde{x}^{\prime} \in S_{x^{\prime}, h} \quad\left|\nabla \widetilde{F}_{0}\left(\tilde{x}^{\prime}\right)\right|^{2} & \leq \frac{C}{h^{5}} \int_{\mathcal{Q}_{0}\left(\tilde{x}^{\prime}, h\right)} \operatorname{dist}^{2}(\nabla y, S O(3)) \\
& \leq \frac{C}{h^{5}} \int_{S_{x^{\prime}, 2 h} \times(0, h)} \operatorname{dist}^{2}(\nabla y, S O(3)), \tag{2.13}
\end{align*}
$$

and henceforth, integrating the pointwise estimate in (2.13) yields

$$
\forall \tilde{x}^{\prime} \in S_{x^{\prime}, h} \quad\left|\widetilde{F}_{0}\left(\tilde{x}^{\prime}\right)-\widetilde{F}_{0}\left(x^{\prime}\right)\right|^{2} \leq \frac{C}{h^{3}} \int_{S_{x^{\prime}, 2 h} \times(0, h)} \operatorname{dist}^{2}(\nabla y, S O(3)),
$$

which combined with (2.12) gives

$$
\begin{equation*}
\forall \tilde{x}^{\prime} \in S_{x^{\prime}, h} \quad\left|\widetilde{F}_{0}\left(\tilde{x}^{\prime}\right)-R_{x^{\prime}}\right|^{2} \leq \frac{C}{h^{3}} \int_{S_{x^{\prime}, 2 h} \times(0, h)} \operatorname{dist}^{2}(\nabla y, S O(3)) . \tag{2.14}
\end{equation*}
$$

On the other hand, applying the triangle inequality we have

$$
\forall z \in \mathcal{Q}_{0}\left(x^{\prime}, h\right) \quad\left|\widetilde{F}_{0}\left(z^{\prime}\right)-\nabla y(z)\right| \leq\left|\widetilde{F}_{0}\left(z^{\prime}\right)-R_{x^{\prime}}\right|+\left|R_{x^{\prime}}-\nabla y(z)\right|
$$

which leads, in view of (2.11) and (2.14), to

$$
\int_{\mathcal{Q}_{0}\left(x^{\prime}, h\right)}\left|\widetilde{F}_{0}\left(z^{\prime}\right)-\nabla y(z)\right|^{2} d z \leq C \int_{S_{x^{\prime}, 2 h} \times(0, h)} \operatorname{dist}^{2}(\nabla y, S O(3)) .
$$

We cover $E_{0}$ with a lattice of non-overlapping squares $S_{x_{i}^{\prime}, h}$ such that $S_{x_{i}^{\prime}, 2 h} \subset U_{0}$ for all $i$. Summing the last estimate over $i$ we obtain

$$
\begin{equation*}
\frac{1}{h} \int_{E_{0} \times(0, h)}\left|\widetilde{F}_{0}\left(z^{\prime}\right)-\nabla y(z)\right|^{2} d z \leq \frac{C}{h} \int_{U_{0} \times(0, h)} \operatorname{dist}^{2}(\nabla y, S O(3)) . \tag{2.15}
\end{equation*}
$$

Similarly, intergrating (2.13) over the cubes $\mathcal{Q}_{0}\left(x_{i}^{\prime}, h\right)$ and summing up over $i$ yields

$$
\begin{equation*}
\int_{E_{0}}\left|\nabla \widetilde{F}_{0}\left(z^{\prime}\right)\right|^{2} d z^{\prime}=\frac{1}{h} \int_{E_{0} \times(0, h)}\left|\nabla \widetilde{F}_{0}\left(z^{\prime}\right)\right|^{2} d z \leq \frac{C}{h^{3}} \int_{U_{0} \times(0, h)} \operatorname{dist}^{2}(\nabla y, S O(3)) . \tag{2.16}
\end{equation*}
$$

For further use we establish a local version of the $L^{2}$ estimate in (2.10). For $x^{\prime} \in E_{0}$ we define

$$
\tilde{d}_{0}\left(x^{\prime}\right):= \begin{cases}\left|\widetilde{F}_{0}\left(x^{\prime}\right)-\mathrm{Id}\right| & \text { if } x^{\prime} \in K_{c_{0} h} \\ \operatorname{dist}\left(\widetilde{F}_{0}\left(x^{\prime}\right), S O(3)\right) & \text { otherwise }\end{cases}
$$

We cover $E_{0}$ as before by the non-overlapping squares $S_{x_{i}^{\prime}, h}$. If $K_{c_{0} h} \cap \overline{S_{x_{i}^{\prime}, h}} \neq \emptyset$, then $R_{x_{i}^{\prime}}=\mathrm{Id}$ and so integrating (2.14) on $S_{x_{i}^{\prime}, h}$ gives

$$
\int_{S_{x_{i}^{\prime}, h}} d_{0}^{2}\left(z^{\prime}\right) d x^{\prime} \leq \int_{S_{x_{i}^{\prime}, h}}\left|\widetilde{F}\left(z^{\prime}\right)-\operatorname{Id}\right|^{2} d z^{\prime} \leq \frac{C}{h} \int_{S_{x^{\prime}, 2 h} \times(0, h)} \operatorname{dist}^{2}(\nabla y, S O(3)) .
$$

Otherwise, if $K_{c_{0} h} \cap \overline{S_{x_{i}^{\prime}, h}}=\emptyset$ we have also by definition of $d_{0}$ and (2.14)

$$
\begin{aligned}
\int_{S_{x_{i}^{\prime}, h}} d_{0}^{2}\left(z^{\prime}\right) d z^{\prime} & =\int_{S_{x_{i}^{\prime}, h}} \operatorname{dist}^{2}\left(\widetilde{F}_{0}\left(z^{\prime}\right), S O(3)\right) d z^{\prime} \\
& \leq \int_{S_{x_{i}^{\prime}, h}}\left|\widetilde{F}\left(z^{\prime}\right)-R_{x^{\prime}}\right|^{2} d z^{\prime} \leq \frac{C}{h} \int_{S_{x^{\prime}, 2 h} \times(0, h)} \operatorname{dist}^{2}(\nabla y, S O(3)) .
\end{aligned}
$$

Summing up the last two inequalities over $i$ we obtain

$$
\begin{equation*}
\left\|\tilde{d}_{0}\right\|_{L^{2}\left(E_{0}\right)}^{2} \leq \frac{C}{h} \int_{U_{0} \times(0, h)} \operatorname{dist}^{2}(\nabla y, S O(3)) . \tag{2.17}
\end{equation*}
$$

Step 2. Boundary estimates: We will construct a matrix field $\widetilde{F}_{j}$ on $E_{j} \cap \Omega$ with useful local estimates. One again, letting $\xi^{\prime}=\Phi\left(x^{\prime}\right)$, we note that

$$
\widetilde{F}_{j}\left(x^{\prime}\right):=\phi_{h} *\left(\bar{F} \circ \Phi_{j}^{-1}\right)\left(\xi^{\prime}\right)
$$

is well-defined for all $x^{\prime} \in E_{j} \cap \Omega$ and $h$ small enough, as the square $S_{\xi^{\prime}, h}$ lies entirely within the open set $\Phi_{j}\left(U_{j} \cap \Omega\right)$ in the upper half-plane.

In this step, we apply Theorem 2.1 (respectively Corollary 2.3(i)) to the Lipschitz domains $\mathcal{Q}_{j}\left(x^{\prime}, h\right):=\Phi_{j}^{-1}\left(S_{\xi^{\prime}, h}\right) \times(0, h)$, noting that the constants in Theorem 2.1 and Corollary 2.3(i)) are invariant under bi-Lipschitz transformations of domains and of boundary conditions under the same transformations. Hence we have

$$
\begin{equation*}
\left\|\nabla y-R_{x^{\prime}}\right\|_{L^{2}\left(\mathcal{Q}_{j}\left(x^{\prime}, h\right)\right)}^{2} \leq C \int_{\mathcal{Q}_{j}\left(x^{\prime}, h\right)} \operatorname{dist}^{2}(\nabla y, S O(3)) \tag{2.18}
\end{equation*}
$$

where $R_{x^{\prime}}=$ Id whenever $K_{c_{j} h} \cap \overline{\Phi_{j}^{-1}\left(S_{\xi^{\prime}, h}\right)} \neq \emptyset$ for a suitable $0<c_{j}<1$. Following in the footsteps of [15, Theorem 6, Step 2] the following estimates similar as in Step 1 are achieved:

$$
\begin{equation*}
\frac{1}{h} \int_{\left(E_{j} \cap \Omega\right)_{\times(0, h)}}\left|\widetilde{F}_{j}\left(z^{\prime}\right)-\nabla y(z)\right|^{2} d z \leq \frac{C}{h} \int_{\left(U_{j} \cap \Omega\right) \times(0, h)} \operatorname{dist}^{2}(\nabla y, S O(3)) \tag{2.20}
\end{equation*}
$$

and

$$
\begin{equation*}
\int_{E_{j} \cap \Omega}\left|\nabla \widetilde{F}_{j}\left(z^{\prime}\right)\right|^{2} d z^{\prime} \leq \frac{C}{h^{3}} \int_{\left(U_{j} \cap \Omega\right) \times(0, h)} \operatorname{dist}^{2}(\nabla y, S O(3)) \tag{2.21}
\end{equation*}
$$

Also if for $x^{\prime} \in E_{j} \cap \Omega$

$$
\tilde{d}_{j}\left(x^{\prime}\right):= \begin{cases}\left|\widetilde{F}_{j}\left(x^{\prime}\right)-\mathrm{Id}\right| & \text { if } x^{\prime} \in K_{c h} \\ \operatorname{dist}\left(\widetilde{F}_{j}\left(x^{\prime}\right), S O(3)\right) & \text { otherwise }\end{cases}
$$

we can prove as before for $d_{0}$ in (2.17):

$$
\begin{equation*}
\left\|\tilde{d}_{j}\right\|_{L^{2}\left(E_{j} \cap \Omega\right)}^{2} \leq \frac{C}{h} \int_{\left(U_{j} \cap \Omega\right) \times(0, h)} \operatorname{dist}^{2}(\nabla y, S O(3)) \tag{2.22}
\end{equation*}
$$

Step 3. Gluing the interior and boundary estimates together:
We now set $\widetilde{F}:=\sum_{j=1}^{N} \theta_{j} \widetilde{F}_{j}$, which is well-defined on $\Omega$. We trivially extend $\theta_{j}, \widetilde{F}$ to $\Omega^{h}$. We note that $\sum \nabla \theta_{j}=0$, and hence

$$
\nabla y-\widetilde{F}=\sum_{j} \theta_{j}\left(\nabla y-\widetilde{F}_{j}\right), \quad \nabla \widetilde{F}=\sum_{j} \theta_{j} \nabla \widetilde{F}_{j}+\sum_{j} \nabla \theta_{j}\left(\widetilde{F}_{j}-\nabla y\right)
$$

Both estimates in (2.9) immediately follow from (2.15), (2.16), (2.20) and (2.21), with constant $C$ depending on the fixed partition of unity $\left\{\left(U_{j}, \theta_{j}\right)\right\}_{j=0}^{N}$, i.e.only on $\Omega$.

To establish (2.10) if $c=\min \left\{c_{j}, j=0, \ldots, N\right\}$, we obtain by the estimates (2.12) and (2.19)

$$
\forall x^{\prime} \in K_{c h} \quad \tilde{d}\left(x^{\prime}\right)^{2}=\left|\sum_{j} \theta_{j}\left(\widetilde{F}_{j}-\mathrm{Id}\right)\right|^{2} \leq C \sum_{j}\left|\widetilde{F}_{j}-\mathrm{Id}\right|^{2} \leq \frac{C}{h^{3}} \int_{\Omega^{h}} \operatorname{dist}^{2}(\nabla y, S O(3))
$$

in view of the fact that if $x^{\prime} \in K_{c_{j} h}$, then necessarily $K_{c_{j} h} \cap \overline{\Phi_{j}^{-1}\left(S_{\xi^{\prime}, h}\right)} \neq \emptyset\left(K_{c_{0} h} \cap \overline{S_{x^{\prime}, h}} \neq \emptyset\right.$ for $j=0$ ) and thus the rotation $R_{x^{\prime}}=\mathrm{Id}$, as previously established.

Otherwise, for any $x^{\prime} \in \Omega$, let

$$
J\left(x^{\prime}\right):=\left\{j \in\{0, \ldots, N\} ; \theta_{j}\left(x^{\prime}\right) \neq 0\right\}
$$

Note that $j \in J\left(x^{\prime}\right)$ implies $x^{\prime} \in E_{j}$. Since the $\Phi_{j}$ are bi-Lipschitz, there exists a constant $C_{0}>0$ depending on the domain $\Omega$ only such that

$$
\bigcup_{j \in J\left(x^{\prime}\right)} \mathcal{Q}_{j}\left(x^{\prime}, h\right) \subset \ddot{\mathcal{Q}}:=\left(B\left(x^{\prime}, C_{0} h\right) \cap \Omega\right) \times(0, h),
$$

where $B\left(x^{\prime}, r\right)$ is the disk of radius $r$ centered at $x^{\prime}$. Applying once again Theorem 2.1 to $y$ on $\ddot{\mathcal{Q}}$, and noting that the constant $C$ still depends only on the Lipschitz constant of $\partial \Omega$, we have for some rotation $\ddot{R}_{x^{\prime}} \in S O(3)$

$$
\begin{equation*}
\left\|\nabla y-\ddot{R}_{x^{\prime}}\right\|_{L^{2}(\ddot{\mathcal{Q}})}^{2} \leq C \int_{\ddot{\mathcal{Q}}} \operatorname{dist}^{2}(\nabla y, S O(3)) \tag{2.23}
\end{equation*}
$$

We note that, since for all $j \in J\left(x^{\prime}\right), \mathcal{Q}_{j}\left(x^{\prime}, h\right) \subset \ddot{\mathcal{Q}}$, we can use $\ddot{R}_{x^{\prime}}$ instead of the rotations used in (2.12) and (2.19), this time obtaining the new bounds

$$
\forall j \in J\left(x^{\prime}\right) \quad\left|\widetilde{F}_{j}\left(x^{\prime}\right)-\ddot{R}_{x^{\prime}}\right| \leq \frac{C}{h^{3}} \int_{\mathcal{Q}_{j}\left(x^{\prime}, h\right)}\left|\nabla y-R_{x^{\prime}}\right|^{2} \leq \frac{C}{h^{3}} \int_{\ddot{\mathcal{Q}}} \operatorname{dist}^{2}(\nabla y, S O(3)) \leq \frac{C}{h^{2}} E^{h}(y)
$$

which establishes

$$
\operatorname{dist}^{2}\left(\widetilde{F}\left(x^{\prime}\right), S O(3)\right) \leq\left|\widetilde{F}\left(x^{\prime}\right)-R_{x^{\prime}}\right|^{2}=\left|\sum_{j \in J\left(x^{\prime}\right)} \theta_{j}\left(x^{\prime}\right)\left(\widetilde{F}_{j}\left(x^{\prime}\right)-\ddot{R}_{x^{\prime}}\right)\right|^{2} \leq \frac{C}{h^{2}} E^{h}(y)
$$

as required for completing the $L^{\infty}$ estimate for $\tilde{d}$ in (2.10) when $x^{\prime} \notin K_{c h}$.
To complete the proof of (2.10), it remain to prove the $L^{2}$ estimate for $\tilde{d}$. We first define for all $x=\left(x^{\prime}, x_{3}\right) \in \Omega^{h}$

$$
d(x):= \begin{cases}|\nabla y(x)-\mathrm{Id}| & \text { if } x^{\prime} \in K_{c h} \\ \operatorname{dist}(\nabla y(x), S O(3)) & \text { otherwise }\end{cases}
$$

and we note that if $x=\left(x^{\prime}, x_{3}\right) \in\left(E_{j} \cap \Omega\right) \times(0, h)$, we have

$$
d(x) \leq\left|\nabla y(x)-\widetilde{F}_{j}\left(x^{\prime}\right)\right|+\tilde{d}_{j}\left(x^{\prime}\right)
$$

Hence the above $L^{2}$ estimates (2.17) and (2.22) on $\tilde{d}_{j}$ obtained in Steps 1 and 2, alongside (2.15) and (2.20) imply

$$
\frac{1}{h} \int_{\left(E_{j} \cap \Omega\right) \times(0, h)} \nabla^{2} \tilde{d}_{j}(x) d x \leq \frac{C}{h} \int_{\left(U_{j} \cap \Omega\right) \times(0, h)} \operatorname{dist}^{2}(\nabla y(x), S O(3)) d x
$$

Summing over $j$ gives

$$
\frac{1}{h}\|d\|_{L^{2}\left(\Omega^{h}\right)}^{2} \leq C E^{h}(y)
$$

which combined with (2.9) proves the $L^{2}$ estimate in (2.10) in view of the fact that

$$
\forall x \in \Omega \quad \tilde{d}\left(x^{\prime}\right) \leq\left|F\left(x^{\prime}\right)-\nabla y(x)\right|+d(x)
$$

Lemma 2.7. Let $\Omega, \Omega^{h}$, $K$ be as defined above. Assume that $h_{0}>0,0<c<1$ be as in Lemma 2.6 and set $\bar{c}=c / 2$. Then, given $h<h_{0}, y \in \mathcal{A}_{K}^{h}$, there exists a matrix field $\widetilde{F}^{\prime}$ which is equal to the identity matrix Id on $K_{\bar{c} h}$ such that the estimates (2.9) and (2.10) still hold true (possibly with a new constant $C$ ) for $\widetilde{F}^{\prime}$, i.e.

$$
\begin{equation*}
\frac{1}{h}\left\|\nabla y-\widetilde{F}^{\prime}\right\|_{L^{2}\left(\Omega^{h}\right)}^{2} \leq C E^{h}(y), \quad\left\|\nabla \widetilde{F}^{\prime}\right\|_{L^{2}(\Omega)}^{2} \leq \frac{C}{h^{2}} E^{h}(y) \tag{2.24}
\end{equation*}
$$

and

$$
\begin{equation*}
\left\|\operatorname{dist}^{2}(\widetilde{F}, S O(3))\right\|_{L^{\infty}(\Omega)}^{2} \leq \frac{C}{h^{2}} E^{h}(y) \tag{2.25}
\end{equation*}
$$

Proof. We let $\widetilde{F}: \Omega \rightarrow \mathbb{R}^{3 \times 3}$ be as in the statement of Lemma 2.6. We introduce a cut-off function $\psi: \mathbb{R}^{2} \rightarrow \mathbb{R}$

$$
\forall x \in \mathbb{R}^{2} \quad \psi(x)= \begin{cases}\frac{\operatorname{dist}\left(x, K_{\bar{c} h}\right)}{\bar{c} h} & \text { if } \operatorname{dist}(x, K) \leq c h \\ 1 & \text { otherwise }\end{cases}
$$

It can be shown that $0 \leq \psi \leq 1, \operatorname{supp} \psi=\overline{K_{\bar{c} h}}, \operatorname{supp}(1-\psi)=\mathbb{R}^{2} \backslash K_{c h}$, and that $\psi$ is a Lipschitz function with

$$
\begin{equation*}
\|\nabla \psi\|_{\infty} \leq \frac{1}{\bar{c} h} \tag{2.26}
\end{equation*}
$$

We now define

$$
\widetilde{F}^{\prime}=\psi \widetilde{F}+(1-\psi) \mathrm{Id}
$$

Note that $\widetilde{F}^{\prime}=\mathrm{Id}$ on $K_{\bar{c} h}$ as required. Also $\widetilde{F}^{\prime}=\widetilde{F}$ on $\Omega \backslash K_{c h}$. We have

$$
\left|\widetilde{F}^{\prime}-\widetilde{F}\right|=\mid \psi \widetilde{F}+(1-\psi) \operatorname{Id}-(\psi \widetilde{F}+(1-\psi) \widetilde{F}|=|(1-\psi)(\widetilde{F}-\mathrm{Id})| \leq \tilde{d}
$$

Hence

$$
\operatorname{dist}\left(\widetilde{F}^{\prime}, S O(3)\right) \leq\left|\widetilde{F}^{\prime}-\widetilde{F}\right|+\operatorname{dist}(\widetilde{F}, S O(3)) \leq 2 \tilde{d} \leq \frac{C}{h^{2}} E^{h}(y)
$$

as required for (2.25). Combining the first estimates in (2.9) and (2.10) with $\left|\widetilde{F}^{\prime}-\widetilde{F}\right| \leq \tilde{d}$ we also obtain

$$
\frac{1}{h}\left\|\nabla y-\widetilde{F}^{\prime}\right\|_{L^{2}\left(\Omega^{h}\right)} \leq C E^{h}(y)
$$

Now we have

$$
\nabla \widetilde{F}^{\prime}=\nabla \psi \otimes(\widetilde{F}-\mathrm{Id})+\psi \nabla \widetilde{F}
$$

which yields, in view of $(2.26)$ and the fact that $\nabla \psi=0$ on the complement of $K_{c h}$,

$$
\left|\nabla \widetilde{F}^{\prime}\right| \leq \frac{C}{h} \tilde{d}+|\nabla \widetilde{F}|
$$

Once again, the $L^{2}$ bounds on $\nabla \widetilde{F}$ and $\widetilde{d}$ in (2.9) and (2.10) imply

$$
\frac{1}{h}\left\|\nabla \widetilde{F}^{\prime}\right\|_{L^{2}\left(\Omega^{h}\right)} \leq \frac{C}{h^{2}} E^{h}(y)
$$

completing the proof of (2.24).
Remark 2.8. In general, by applying a projection on a large ball containing $S O(3)$, we can also assume that $\left|\widetilde{F}^{\prime}\right| \leq C$. By (2.25), this will not be needed under the assumption $E^{h} \leq C h^{2}$.

In the following statement we now arrange $\tilde{F}^{\prime}$ in such a way that it takes values only into $S O(3)$.

Corollary 2.9. Let $\Omega, \Omega^{h}, K$ be as defined above. Then there exist constants $h_{0}>0, C>0$, $0<\bar{c}<1, \delta_{0}>0$, depending only on $\Omega$ and $K$, such that if for any $h<h_{0}, y \in \mathcal{A}_{K}^{h}$ we have $E^{h}(y) \leq \delta_{0} h^{2}$, then there exists a matrix valued mapping $R \in W^{1,2}(\Omega, S O(3))$ such that $\left.R\right|_{K_{\bar{c} h}} \equiv \mathrm{Id}$ and the following estimates hold true:

$$
\begin{equation*}
\frac{1}{h}\|\nabla y-R\|_{L^{2}\left(\Omega^{h}\right)}^{2} \leq C E^{h}(y), \quad\|\nabla R\|_{L^{2}(\Omega)}^{2} \leq \frac{C}{h^{2}} E^{h}(y) \tag{2.27}
\end{equation*}
$$

Proof. We let $\widetilde{F}, h_{0}, \bar{c}$ be chosen according to Lemma 2.24. The result follows similarly as in $[15$, Remark 5]: Under the assumptions of the Corollary $2.9,(2.25)$ implies that $\widetilde{F}^{\prime}$ is in a $\left(C \delta_{0}\right)^{\frac{1}{2}}$ tubular neighborhood $\mathcal{U}$ of $S O(3)$. For $\delta_{0}$ small enough, the orthogonal projection $\pi$ from $\mathcal{U}$ onto $S O(3)$ is well-defined and Lipschitz. We let $R:=\pi \circ \widetilde{F}^{\prime}$, and note that

$$
\begin{aligned}
|\nabla y-R| & \leq\left|\nabla y-\widetilde{F}^{\prime}\right|+\left|\widetilde{F}^{\prime}-R\right| \\
& =\left|\nabla y-\widetilde{F}^{\prime}\right|+\operatorname{dist}\left(\widetilde{F}^{\prime}, S O(3)\right) \\
& \leq 2\left|\nabla y-\widetilde{F}^{\prime}\right|+\operatorname{dist}(\nabla y, S O(3))
\end{aligned}
$$

and that

$$
|\nabla R| \leq\|\nabla \pi\|_{L^{\infty}}\left|\nabla \tilde{F}^{\prime}\right|
$$

which combined with (2.24) imply together (2.27).
Let for $x_{3} \in(0,1), \tilde{y}^{h}\left(x^{\prime}, x_{3}\right):=y^{h}\left(x^{\prime}, h x_{3}\right)=\left(\left(y^{h}\right)^{\prime}\left(x^{\prime}, h x_{3}\right), y_{3}^{h}\left(x^{\prime}, h x_{3}\right)\right)$. We consider the mappings

$$
\operatorname{id}^{h}: \Omega^{h} \rightarrow \mathbb{R}^{3}, \quad \operatorname{id}^{h}\left(x^{\prime}, x_{3}\right)=\left(x^{\prime}, x_{3}\right) ; \quad \operatorname{id}: \Omega^{1} \rightarrow \mathbb{R}^{3} \quad \operatorname{id}\left(x^{\prime}, x_{3}\right)=\left(\operatorname{id}^{h}\right)^{\prime}\left(x^{\prime}, x_{3}\right)=\left(x^{\prime}, 0\right)
$$

and the displacement fields

$$
\begin{equation*}
u_{h} \in W^{1,2}(\Omega, \mathbb{R}), \quad U_{h}:=\frac{1}{h} \int_{0}^{h}\left(y_{3}^{h}\left(\cdot, x_{3}\right)-h x_{3}\right) d x_{3}, \quad u_{h}:=\frac{1}{h^{\beta / 2-1}} U_{h} \tag{2.28}
\end{equation*}
$$

$w_{h} \in W^{1,2}\left(\Omega, \mathbb{R}^{2}\right), \quad W_{h}:=\frac{1}{h} \int_{0}^{h}\left(\left(y^{h}\right)^{\prime}-\mathrm{id}\right)\left(\cdot, x_{3}\right) d x_{3}, \quad w_{h}:=\frac{1}{h^{\delta}} W_{h}$ for $\delta:=\min \{\beta-2, \beta / 2\}$.
Theorem 2.10. Let $\Omega, \Omega^{h}$, $K$ be as above, and assume that $\operatorname{Cap}_{1,2}(K)>0$ and $\beta>2$. Assume that for a sequence as $h \rightarrow 0, y^{h} \in \mathcal{A}_{K}^{h}, E^{h}\left(y^{h}\right) \precsim h^{\beta}$. Then, up to a subsequence as $h \rightarrow 0$, we have
(i) $\frac{1}{h}\left\|y^{h}-\operatorname{id}^{h}\right\|_{W^{1,2}\left(\Omega^{h}\right)}^{2} \longrightarrow 0,\left(\right.$ in particular $\tilde{y}^{h} \longrightarrow \mathrm{id}$ in $\left.W^{1,2}\left(\Omega^{1}\right)\right)$.
(ii) $u_{h} \xrightarrow{\text { strongly in } W^{1,2}} u \in H_{0, K}^{2}(\Omega)$.
(iii) $w_{h} \xrightarrow{\text { weakly in } W^{1,2}} w \in W^{1,2}(\Omega)$.
(iv) If $2<\beta<4$, then $w_{h} \xrightarrow{\text { strongly in } W^{1,2}} w \in H_{0, K}^{1}(\Omega)$; and

$$
\operatorname{sym} \nabla w+\frac{1}{2} \nabla u \otimes \nabla u \equiv 0 \quad \text { in } \Omega
$$

(v) For $\beta \geq 4$, assume that $\operatorname{sym} \nabla w_{h}$ converges strongly in $L^{2}(\Omega)$, and that moreover the family

$$
\mathcal{Y}:=\frac{1}{h^{\beta}}\left(\frac{1}{h} \int_{0}^{h} \operatorname{dist}^{2}\left(\nabla y^{h}\left(\cdot, x_{3}\right), S O(3)\right) d x_{3}\right)
$$

is equi-integrable over $\Omega$, in the sense that for all $\varepsilon>0$, there exists $\delta>0$ such that for all $B \subset \Omega,|B|<\delta$, we have

$$
\frac{1}{h^{\beta}}\left(\frac{1}{h} \int_{B \times(0, h)} \operatorname{dist}^{2}\left(\nabla y^{h}, S O(3)\right) d x\right)<\varepsilon
$$

Then $w_{h} \xrightarrow{\text { strongly in } W^{1,2}} w \in H_{0, K}^{1}(\Omega)$.

Proof. Let $h_{0}, \delta_{0}, \bar{c}, C$ be as in Corollary 2.9. By the assumption $\beta>2$, and if necessary by choosing a smaller value for $h_{0}$, we can assure that $h<h_{0}$ implies $E^{h}\left(y^{h}\right) \leq \delta_{0} h^{2}$ and so Corollary 2.9 applies: For each $y^{h}$ with $h<h_{0}$, we denote the associated rotation field by $R^{h}$, and hence (2.27) implies

$$
\begin{equation*}
\frac{1}{h}\left\|\nabla y^{h}-R^{h}\right\|_{L^{2}\left(\Omega^{h}\right)}^{2} \leq C h^{\beta}, \quad\left\|\nabla\left(R^{h}-\mathrm{Id}\right)\right\|_{L^{2}(\Omega)}^{2} \leq C h^{\beta-2} \tag{2.30}
\end{equation*}
$$

where $R^{h} \equiv$ Id on $K_{\bar{c} h}$. We let

$$
\begin{equation*}
A^{h}:=\frac{1}{h^{\beta / 2-1}}\left(R^{h}-\mathrm{Id}\right) \tag{2.31}
\end{equation*}
$$

By the second estimate in (2.30), $\nabla A^{h}$ is uniformly bounded in $L^{2}(\Omega)$ and that $A^{h} \equiv 0$ on $K_{\bar{c} h}$. Moreover, by the Poincaré inequality proved in Corollary A. 4 the sequence $A^{h}$ satisfies

$$
\left\|A^{h}\right\|_{L^{2}(\Omega)} \leq C\left\|\nabla A^{h}\right\|_{L^{2}(\Omega)}
$$

and hence up to a subsequence, $A^{h}$ converges weakly in $W^{1,2}$ to $A$. Note that by Lemma A.1, we have

$$
\begin{equation*}
A \in H_{0, K}^{1}(\Omega) \tag{2.32}
\end{equation*}
$$

We have moreover

$$
\begin{equation*}
\operatorname{sym} R^{h}-\mathrm{Id}=-\frac{1}{2}\left(R^{h}-\mathrm{Id}\right)^{T}\left(R^{h}-\mathrm{Id}\right) \tag{2.33}
\end{equation*}
$$

which implies that

$$
\operatorname{sym} A^{h}=-\frac{1}{2} h^{\beta / 2-1}\left(A^{h}\right)^{T} A^{h},
$$

and as a consequence, passing to the limit, we obtain that sym $A \equiv 0$ in $\Omega$. Now, using (2.33), this time rescaled with $h^{\beta-2}$, Sobolev embeddings, and the fact that $A^{T} A=-A^{2}$ for any skew-symmetric matrix, we obtain that

$$
\begin{equation*}
\forall p<2 \quad B^{h}:=\frac{1}{h^{\beta-2}}\left(\operatorname{sym} R^{h}-\mathrm{Id}\right)=-\frac{1}{2}\left(A^{h}\right)^{T} A^{h} \rightharpoonup \frac{A^{2}}{2} \text { weakly in } W^{1, p}(\Omega) \tag{2.34}
\end{equation*}
$$

which also implies the strong convergence of $B^{h}$ in $L^{q}(\Omega)$ for all $q<\infty$.
We first note that the convergence of $u_{h}$ and $w_{h}$ in $W^{1,2}$ as we shall prove below in parts (ii) and (iii) imply (i) in a straightforward manner. Indeed, we first remark that by compactness of Sobolev embeddings, $A^{h}$ must converge strongly to $A$ in $L^{2}(\Omega)$. Now the first estimate in (2.30) implies

$$
\begin{align*}
\frac{1}{h}\left\|\frac{1}{h^{\frac{\beta}{2}-1}}\left(\nabla y^{h}-\mathrm{Id}\right)-A\right\|_{L^{2}\left(\Omega^{h}\right)}^{2} & \leq \frac{C}{h}\left(\left\|\frac{1}{h^{\frac{\beta}{2}-1}}\left(\nabla y^{h}-R^{h}\right)\right\|_{L^{2}\left(\Omega^{h}\right)}^{2}+\left\|A^{h}-A\right\|_{L^{2}\left(\Omega^{h}\right)}^{2}\right)  \tag{2.35}\\
& \leq C\left(h^{2}+\left\|A^{h}-A\right\|_{L^{2}(\Omega)}^{2}\right) \xrightarrow{h \rightarrow 0} 0,
\end{align*}
$$

which in particular yields

$$
\begin{equation*}
\frac{1}{h}\left\|\nabla\left(y^{h}-\mathrm{id}^{h}\right)\right\|_{L^{2}\left(\Omega^{h}\right)}^{2} \leq C h^{\beta-2} \xrightarrow{h \rightarrow 0} 0 \tag{2.36}
\end{equation*}
$$

To obtain an $L^{2}$ estimate on $y^{h}-\mathrm{id}^{h}$, for a.e. $x^{\prime} \in \Omega$ we apply the Poincaré inequality on the interval $\left\{x^{\prime}\right\} \times(0, h)$ and integrate over $\Omega$ to obtain

$$
\begin{aligned}
\frac{1}{h}\left\|\left(y^{h}\right)^{\prime}-\left(\mathrm{id}^{h}\right)^{\prime}-W_{h}\right\|_{L^{2}\left(\Omega^{h}\right)}^{2} & =\frac{1}{h}\left\|\left(y^{h}\right)^{\prime}-\mathrm{id}-f_{0}^{h}\left(\left(y^{h}\right)^{\prime}-\mathrm{id}\right)\left(\cdot, x_{3}\right) d x_{3}\right\|_{L^{2}\left(\Omega^{h}\right)}^{2} \\
& \leq C h^{2}\left(\frac{1}{h}\left\|\partial_{3}\left(\left(y^{h}\right)^{\prime}-\left(\mathrm{id}^{h}\right)^{\prime}\right)\right\|_{L^{2}(\Omega)}^{2}\right) \\
& \leq C h^{2}\left(\frac{1}{h}\left\|\nabla\left(y^{h}-\mathrm{id}^{h}\right)\right\|_{L^{2}(\Omega)}^{2}\right)
\end{aligned}
$$

for the first two components of $y^{h}-\mathrm{id}^{h}$, and

$$
\begin{aligned}
\frac{1}{h}\left\|y_{3}^{h}-\mathrm{id}_{3}^{h}-U_{h}\right\|_{L^{2}\left(\Omega^{h}\right)}^{2} & =\frac{1}{h}\left\|y_{3}^{h}-\mathrm{id}_{3}^{h}-f_{0}^{h}\left(y_{3}^{h}-\mathrm{id}_{3}^{h}\right)\left(\cdot, x_{3}\right) d x_{3}\right\|_{L^{2}\left(\Omega^{h}\right)}^{2} \\
& \leq C h^{2}\left(\frac{1}{h}\left\|\partial_{3}\left(y_{3}^{h}-\mathrm{id}_{3}^{h}\right)\right\|_{L^{2}(\Omega)}^{2}\right) \\
& \leq C h^{2}\left(\frac{1}{h}\left\|\nabla\left(y^{h}-\mathrm{id}^{h}\right)\right\|_{L^{2}(\Omega)}^{2}\right)
\end{aligned}
$$

Combining the last two estimates with (2.36) we hence obtain

$$
\begin{aligned}
\frac{1}{h}\left\|y^{h}-\mathrm{id}^{h}\right\|_{L^{2}\left(\Omega^{h}\right)}^{2} & \leq C\left(\left\|U_{h}\right\|_{L^{2}(\Omega)}^{2}+\left\|W_{h}\right\|_{L^{2}(\Omega)}^{2}+h^{\beta}\right) \\
& \leq C\left(h^{\beta-2}\left\|u_{h}\right\|_{L^{2}(\Omega)}^{2}+h^{2 \delta}\left\|w_{h}\right\|_{L^{2}(\Omega)}^{2}+h^{\beta}\right) \xrightarrow{h \rightarrow 0} 0
\end{aligned}
$$

where we used (ii) and (iii).
To see (ii), first observe that for $\nabla^{\prime}=\nabla_{x^{\prime}}$

$$
\begin{aligned}
\left\|\nabla u_{h}-\left(A_{31}, A_{32}\right)\right\|_{L^{2}(\Omega)} & =\left\|\frac{1}{h^{\beta / 2-1}} \nabla^{\prime} \int_{0}^{1} \tilde{y}_{3}^{h}-h x_{3} d x_{3}-\left(A_{31}, A_{32}\right)\right\|_{L^{2}(\Omega)} \\
& =\left\|\int_{0}^{1} \frac{1}{h^{\beta / 2-1}} \nabla^{\prime} \tilde{y}_{3}^{h}-h x_{3} d x_{3}-\left(A_{31}, A_{32}\right)\right\|_{L^{2}(\Omega)} \\
& \leq \frac{1}{\sqrt{h}}\left\|\frac{1}{h^{\beta / 2-1}} \nabla\left(y^{h}-\mathrm{id}^{h}\right)-A\right\|_{L^{2}\left(\Omega^{h}\right)} \\
& =\frac{1}{\sqrt{h}}\left\|\frac{1}{h^{\beta / 2-1}}\left(\nabla y^{h}-\mathrm{Id}\right)-A\right\|_{L^{2}\left(\Omega^{h}\right)} \\
& \leq \frac{1}{\sqrt{h}}\left\|\frac{1}{h^{\beta / 2-1}}\left(\nabla y^{h}-R^{h}\right)\right\|_{L^{2}\left(\Omega^{h}\right)}+\left\|\frac{1}{h^{\beta / 2-1}}\left(R^{h}-\mathrm{Id}\right)-A\right\|_{L^{2}(\Omega)} \\
& \leq C \sqrt{h}+\left\|A^{h}-A\right\|_{L^{2}(\Omega)} \xrightarrow{h \rightarrow 0} 0
\end{aligned}
$$

up to a subsequence, where we used (2.30) and the convergence of $A^{h}$ to $A$ in the last line. Therefore for $a_{h}:=f_{\Omega} u_{h}$ we obtain by the Poincaré-Sobolev inequalities that for some $\tilde{u} \in$ $W^{2,2}(\Omega)$, with $\nabla \tilde{u}=\left(A_{31}, A_{32}\right)$,

$$
\left\|u_{h}-a_{h}-\tilde{u}\right\|_{W^{1,2}(\Omega)} \xrightarrow{h \rightarrow 0} 0
$$

(To see this, we first establish the weak convergence $u_{h}-a_{h}$ and then use the strong convergence of $\nabla u_{h}$.) Note that by Sobolev embedding theorems $\tilde{u} \in C^{0}(\bar{\Omega})$, and hence the value of $\tilde{u}$ for all points on $K$ is well-defined. By Lemma 2.11 proven below, we have for q.e. $x \in K$

$$
\lim _{h \rightarrow 0} f_{B(x, \bar{c} h) \cap \Omega}\left|u_{h}-a_{h}-\tilde{u}\right|=0 \Longrightarrow \lim _{h \rightarrow 0} f_{B(x, \bar{c} h) \cap \Omega}\left(u_{h}-a_{h}-\tilde{u}\right)=0
$$

which implies

$$
\begin{equation*}
\lim _{h \rightarrow 0} f_{B(x, \bar{c} h) \cap \Omega} u_{h}-a_{h}=\lim _{h \rightarrow 0} f_{B(x, \bar{c} h) \cap \Omega} \tilde{u}=\tilde{u}(x) . \tag{2.37}
\end{equation*}
$$

On the other hand, since $y^{h}=\mathrm{id}$ on $K_{h}$ :

$$
\begin{aligned}
\int_{B(x, \bar{c} h) \cap \Omega}\left|U_{h}\right|^{2} & \leq \frac{1}{h} \int_{(B(x, \bar{c} h) \cap \Omega) \times(0, h)}\left|y_{3}^{h}-x_{3}\right|^{2} \\
& \leq C h^{2}\left(\frac{1}{h} \int_{(B(x, \bar{c} h) \cap \Omega) \times(0, h)}\left|\partial_{3}\left(y_{3}^{h}-x_{3}\right)\right|^{2}\right) \\
& \leq C h^{2}\left(\frac{1}{h} \int_{(B(x, \bar{c} h) \cap \Omega) \times(0, h)}\left|\nabla y^{h}-\mathrm{Id}\right|^{2}\right) \\
& \leq C h^{2}\left(\frac{1}{h} \int_{(B(x, \bar{c}) \cap \Omega) \times(0, h)}\left|\nabla y^{h}-R^{h}\right|^{2}\right) \leq C h^{\beta+2},
\end{aligned}
$$

since $R^{h}=$ Id on $K_{\bar{c} h}$. Therefore, since $|B(x, \bar{c} h) \cap \Omega| \geq C(\Omega) h^{2}$ we obtain

$$
f_{B(x, \bar{c} h) \cap \Omega}\left|u_{h}\right|^{2} \leq C h^{2},
$$

which implies

$$
\lim _{h \rightarrow 0} f_{B(x, \bar{c} h) \cap \Omega}\left|u_{h}\right|=0
$$

yielding, combined with (2.37), that $a=\lim _{h \rightarrow 0} a_{h} \in \mathbb{R}$ exists, and that for q.e. $x \in K, \tilde{u}(x)=-a$. Therefore for $u:=\tilde{u}+a$ :

$$
u_{h}=\left(u_{h}-a_{h}\right)+a_{h} \xrightarrow{\text { in } W^{1,2}} \tilde{u}+a=u \in W^{2,2}(\Omega),
$$

and the continuous representative of $u$ vanishes on $K$. More precisely, we obtain by Lemma 2.11, since $u_{h}-u$ converges strongly to 0 in $W^{1,2}(\Omega)$ :

$$
f_{B(x, \overline{c h}) \cap \Omega}|u| \leq f_{B(x, \overline{c h}) \cap \Omega}\left|u-u_{h}\right|+f_{B(x, \overline{c h}) \cap \Omega}\left|u_{h}\right| \xrightarrow{h \rightarrow 0} 0 .
$$

Therefore, by definition, $u \in H_{0, K}^{1}(\Omega)$, which alongside the fact that by (2.32) $\nabla u=\left(A_{31}, A_{32}\right) \in$ $H_{0, K}^{1}(\Omega)$, completes the proof of (ii).
To establish (iii) and (iv) we proceed in the same manner as in (ii), with a caveat. Let $\overparen{B}$ be the $2 \times 2$ upper-left sub-matrix of $B:=A^{2} / 2$. We have

$$
\begin{aligned}
\| \operatorname{sym} \nabla w_{h}-h^{(\beta-2)-\delta} & \widehat{B}\left\|_{L^{2}(\Omega)}=\right\| \operatorname{sym}\left(\frac{1}{h^{\delta}} \nabla^{\prime} \int_{0}^{1}\left(\tilde{y}^{h}\right)^{\prime}-x^{\prime} d x_{3}\right)-h^{(\beta-2)-\delta} \widehat{B} \|_{L^{2}(\Omega)} \\
& =\left\|\int_{0}^{1} \frac{1}{h^{\delta}} \operatorname{sym}\left(\nabla^{\prime}\left(\left(\tilde{y}^{h}\right)^{\prime}-x^{\prime}\right) d x_{3}\right)-h^{(\beta-2)-\delta} \widehat{B}\right\|_{L^{2}(\Omega)} \\
& \leq \frac{1}{\sqrt{h}}\left\|\frac{1}{h^{\delta}} \operatorname{sym}\left(\nabla\left(y^{h}-\mathrm{id}^{h}\right)\right)-h^{(\beta-2)-\delta} B\right\|_{L^{2}\left(\Omega^{h}\right)} \\
& =\frac{1}{\sqrt{h}}\left\|\frac{1}{h^{\delta}} \operatorname{sym}\left(\nabla y^{h}-\mathrm{Id}\right)-h^{(\beta-2)-\delta} B\right\|_{L^{2}\left(\Omega^{h}\right)} \\
& \leq \frac{1}{\sqrt{h}}\left\|\frac{1}{h^{\delta}} \operatorname{sym}\left(\nabla y^{h}-R^{h}\right)\right\|_{L^{2}\left(\Omega^{h}\right)}+\left\|\frac{1}{h^{\delta}} \operatorname{sym}\left(R^{h}-\mathrm{Id}\right)-h^{(\beta-2)-\delta} B\right\|_{L^{2}(\Omega)} \\
& \leq C h^{(\beta / 2)-\delta}+h^{(\beta-2)-\delta}\left\|B^{h}-B\right\|_{L^{2}(\Omega)} .
\end{aligned}
$$

by (2.30), (2.34). If $2<\beta<4$, then $\delta=\beta-2<\beta / 2$ and we deduce that $\operatorname{sym} \nabla w_{h}$ converges in $L^{2}(\Omega)$ to $\widehat{B}$ by (2.34). Otherwise, if $\beta \geq 4$, then $\delta=\beta / 2$ and $h^{(\beta-2)-\delta} \rightarrow 0$, and as a consequence
$\left\|\operatorname{sym} \nabla w_{h}\right\|_{L^{2}(\Omega)}$ is merely bounded. This implies, through applying Korn's inequality on $\Omega$, that if for all $x \in \Omega$ we set $\tilde{w}_{h}(x):=w_{h}(x)-D^{h} x-b_{h}$, where

$$
b_{h}:=f_{\Omega} w_{h} \quad \text { and } \quad D^{h}=\left[\begin{array}{cc}
0 & d_{h} \\
-d_{h} & 0
\end{array}\right]:=f_{\Omega} \operatorname{skew}\left(\nabla w_{h}\right)
$$

there exists $\tilde{w} \in W^{1,2}\left(\Omega, \mathbb{R}^{2}\right)$, with $\operatorname{sym} \nabla \tilde{w}=\overparen{B}$ if $2<\beta<4$, so that up to a subsequence

$$
\begin{array}{ll}
\text { if } 2<\beta<4 & \tilde{w}_{h} \xrightarrow{\text { in } W^{1,2}} \tilde{w} \\
\text { if } \beta \geq 4 & \tilde{w}_{h} \xrightarrow{\text { weakly in } W^{1,2}} \tilde{w} .
\end{array}
$$

If $2<\beta<4$ we obtain through Lemma 2.11 that for q.e. $x \in \bar{\Omega}$

$$
\lim _{h \rightarrow 0} f_{B(x, \bar{c} h) \cap \Omega} \tilde{w}_{h}=\lim _{h \rightarrow 0} f_{B(x, \bar{c} h) \cap \Omega} \tilde{w}=\tilde{w}^{*}(x)
$$

by applying an extension operator of Appendix A and [13, Theorem 4.8.1], where $\tilde{w}^{*}(x)$ is the precise representative of $\mathcal{E}(\tilde{w})$. We now observe that since $y^{h}=$ id on $K_{h}$ and $R^{h} \equiv \operatorname{Id}$ on $K_{\bar{c} h}$ :

$$
\begin{aligned}
\int_{B(x, \bar{c} h) \cap \Omega}\left|W_{h}\right|^{2} & \leq \frac{1}{h} \int_{(B(x, \bar{c} h) \cap \Omega) \times(0, h)}\left|\left(y^{h}\right)^{\prime}(x)-x^{\prime}\right|^{2} \\
& \leq C h^{2}\left(\frac{1}{h} \int_{(B(x, \bar{c} h) \cap \Omega) \times(0, h)}\left|\partial_{3}\left(\left(y^{h}\right)^{\prime}-x^{\prime}\right)\right|^{2}\right) \\
& \leq C h^{2}\left(\frac{1}{h} \int_{(B(x, \bar{c} h) \cap \Omega) \times(0, h)}\left|\nabla y^{h}-\mathrm{Id}\right|^{2}\right) \\
& \leq C h^{2}\left(\frac{1}{h} \int_{(B(x, \bar{c} h) \cap \Omega) \times(0, h)}\left|\nabla y^{h}-R^{h}\right|^{2}\right) \leq C h^{\beta+2}
\end{aligned}
$$

Thus, once again through $|B(x, \bar{c} h) \cap \Omega| \geq C(\Omega) h^{2}$ we obtain

$$
f_{B(x, \bar{c} h) \cap \Omega}\left|w_{h}\right|^{2} \leq C h^{\beta-2 \delta}
$$

which implies for $2<\beta<4$ :

$$
\lim _{h \rightarrow 0} f_{B(x, \bar{c} h) \cap \Omega}\left|w_{h}\right|=0
$$

Therefore we obtain for q.e. $x \in K$

$$
\lim _{h \rightarrow 0}\left(d_{h} f_{B(x, \bar{c} h) \cap \Omega} y^{\perp} d y+b_{h}\right)=\lim _{h \rightarrow 0}\left(D^{h} f_{B(x, \bar{c} h) \cap \Omega} y^{\prime} d y+b_{h}\right)=-\tilde{w}^{*}(x)
$$

But by regularity of the mapping $y \rightarrow y^{\perp}$, for all $x \in \bar{\Omega}$

$$
\lim _{h \rightarrow 0} f_{B(x, \bar{c} h) \cap \Omega} y^{\perp} d y=x^{\perp}
$$

so we conclude with

$$
\text { for q.e. } x \in K \quad \lim _{h \rightarrow 0}\left(d_{h} x^{\perp}+b_{h}\right)=-\tilde{w}^{*}(x) \text {. }
$$

Note that the q.e. existence of the above limit implies its existence for at least two distinct values of $x \in K$, from which immediately follows that $d:=\lim _{h \rightarrow 0} d_{h}$ and $b:=\lim _{h \rightarrow 0} b_{h}$ exist and subsequently that for q.e. $x \in K$,

$$
\tilde{w}^{*}(x)=-\left(d x^{\perp}+b\right)
$$

Now, $w_{h}-\tilde{w}_{h}(x)=D^{h} x+b_{h} \rightarrow d x^{\perp}+b$ which yields that

$$
w_{h}=\left(w_{h}-\tilde{w}_{h}\right)+\tilde{w}_{h} \xrightarrow{i n W^{1,2}} w:=\tilde{w}+d x^{\perp}+b .
$$

Applying Lemma 2.11 once again in an argument similar to what we presented for $u$ in part (ii), yields $w \in H_{0, K}^{1}(\Omega)$.

Finally, remember that $\overparen{F}$ is the $2 \times 2$ upper-left sub-matrix of $F \in R^{3 \times 3}$. First we observe that by (2.35)

$$
\begin{aligned}
\left\|\frac{h^{\delta}}{h^{\beta / 2-1}} \nabla w_{h}-\overparen{A}\right\|_{L^{2}(\Omega)}^{2} & =\left\|\frac{1}{h^{\beta / 2-1}}\left(f_{0}^{h} \nabla^{\prime}\left(\left(y^{h}\right)^{\prime}-\left(\mathrm{id}^{h}\right)^{\prime}\right)\left(\cdot, x_{3}\right) d x_{3}\right)-\overparen{A}\right\|_{L^{2}(\Omega)}^{2} \\
& =\|f_{0}^{h}(\overbrace{\frac{1}{h^{\beta / 2-1}}\left(\nabla y^{h}-\mathrm{Id}\right)-A}^{)}) d x_{3}\|_{L^{2}(\Omega)}^{2} \\
& \leq \frac{1}{h}\left(\left\|\frac{1}{h^{\beta / 2-1}}\left(\nabla y^{h}-\mathrm{Id}\right)-A\right\|_{L^{2}\left(\Omega^{h}\right)}^{2}\right) \xrightarrow{h \rightarrow 0} 0 .
\end{aligned}
$$

But since $\beta>2$, we have $\delta>\beta / 2-1$, and hence, in view of the boundedness of $\left\|\nabla w_{h}\right\|_{L^{2}}$ and by passing to the limit in $h \rightarrow 0$ we obtain $\overparen{A}=0$. Remember that $A^{T}=-A$, and that $\nabla u=\left(A_{31}, A_{32}\right)$. Hence, overall, straightforward calculation gives

$$
A=e_{3} \otimes \nabla u-\nabla u \otimes e_{3}=\left[\begin{array}{ccc}
0 & 0 & -\partial_{1} u  \tag{2.38}\\
0 & 0 & -\partial_{2} u \\
\partial_{1} u & \partial_{2} u & 0
\end{array}\right]
$$

which implies

$$
\begin{equation*}
\widehat{B}=\overparen{\overbrace{\frac{A^{2}}{2}}^{2}}=-\frac{1}{2} \nabla u \otimes \nabla u \tag{2.39}
\end{equation*}
$$

However, we already know that if $2<\beta<4$, then $\operatorname{sym} \nabla w=\widehat{B}$. Hence the proof of (iv) is complete.

In order to see (v), it is sufficient to repeat the same argument as in (iv) using the new assumptions. The weak convergence of $\operatorname{sym} \nabla w_{h}$ being replaced by strong convergence, we deduce the strong convergence of the adjusted sequence $\tilde{w}_{h}$. Meanwhile, a careful reexamination of the proof of Corollary 2.9 alongside the observations made in [15, Proposition 4] (see also [11, Corollary 4.2]) implies that the family

$$
\mathscr{Y}=\frac{1}{h^{\beta}}\left(\frac{1}{h} \int_{0}^{h}\left|\nabla y^{h}\left(\cdot, x_{3}\right)-R^{h}\right|^{2} d x_{3}\right)
$$

is equi-integrable over $\Omega$. Now, since $\beta=2 \delta$ for $\beta \geq 4$, the estimate

$$
\int_{B(x, \bar{c} h) \cap \Omega}\left|W_{h}\right|^{2} \leq C h^{2}\left(\frac{1}{h} \int_{(B(x, \bar{c} h) \cap \Omega) \times(0, h)}\left|\nabla y^{h}-R^{h}\right|^{2}\right)
$$

in combination with the equi-integrability of $\mathscr{Y}$ implies once again that

$$
f_{B(x, \bar{c} h) \cap \Omega}\left|w_{h}\right|^{2} \xrightarrow{h \rightarrow 0} 0 .
$$

The conclusion then follows similarly as in (iv).

Lemma 2.11. Let $\Omega \subset \mathbb{R}^{2}$ be a bounded Lipschitz domain, $\bar{c}>0$ and assume that for $a$ subsequence $h_{j} \rightarrow 0, u_{h} \in W^{1,2}(\Omega)$ converge strongly to 0 in $W^{1,2}$. Then, up to a subsequence, and for q.e. $x \in \bar{\Omega}$

$$
\lim _{h \rightarrow 0} f_{B(x, \bar{c} h) \cap \Omega}\left|u_{h}(y)\right| d y=0
$$

Proof. Since $\left|u_{h}\right|$ also converge strongly to $|u|$, without loss of generality we can assume that $u_{h} \geq 0$ a.e. in $\Omega$. We first extend the $u_{h}$ to the whole $\mathbb{R}^{2}$ using the bounded linear extension operator $\mathcal{E}$ introduced in Appendix A, formula (A.6). It directly follows from Proposition A. 2 that

$$
\lim _{h \rightarrow 0} f_{B(x, \bar{c} h)} \mathcal{E}\left(u_{h}\right)(y) d y=0 \Rightarrow \lim _{h \rightarrow 0} f_{B(x, \bar{c} h) \cap \Omega} u_{h}(y) d y=0
$$

Therefore it is sufficient to prove the statement for $\Omega=\mathbb{R}^{2}$. In that case, we choose a suitable subsequence of $u_{h_{j}}$ which satisfies, after relabeling,

$$
\int_{\mathbb{R}^{2}}\left|D u_{h_{j}}\right|^{2} \leq \frac{1}{2^{3 j}}
$$

and we define

$$
B_{j}:=\left\{x \in \mathbb{R}^{2} ; f_{B\left(x, \bar{c} h_{j}\right)} u_{h_{j}}(y) d y>\frac{1}{2^{j}}\right\} \subset\left\{x \in \mathbb{R}^{2} ; f_{B(x, r)} u_{h_{j}}(y) d y>\frac{1}{2^{j}} \text { for some } r>0\right\}
$$

Applying [13, Lemma 4.8.1] to $u_{h_{j}}$, we obtain that

$$
\operatorname{Cap}_{1,2}\left(B_{j}\right) \leq C 2^{2 j} \int_{\mathbb{R}^{2}}\left|D u_{h_{j}}\right|^{2} \leq \frac{C}{2^{j}}
$$

Let

$$
E_{k}:=\bigcup_{j=k}^{\infty} B_{j}
$$

Then

$$
\operatorname{Cap}_{1,2}\left(E_{k}\right) \leq \frac{C}{2^{k-1}}
$$

and for all $x \in \mathbb{R}^{2} \backslash E_{k}$ we have for all $j \geq k$

$$
f_{B\left(x, \bar{c} h_{j}\right)} u_{h_{j}}(y) \leq \frac{1}{2^{3 j}}
$$

Thus by letting

$$
E:=\bigcap_{k=1}^{\infty} E_{k}
$$

we have $\operatorname{Cap}_{1,2}(E)=0$ and

$$
\forall x \in \mathbb{R}^{2} \backslash E \quad \lim _{j \rightarrow \infty} f_{B\left(x, \bar{c} h_{j}\right)} u_{h_{j}}(y)=0
$$

The proof is complete.
Corollary 2.12. Let $\Omega, \Omega^{h}, K, \beta$ be as above. Assume that for a sequence as $h \rightarrow 0, y^{h} \in \mathcal{A}_{K}^{h}$, $E^{h}\left(y^{h}\right) \precsim h^{\beta}$, $R^{h}$ is chosen as in (2.30), and $u$, $w$ are as in Theorem 2.10. Let

$$
G^{h}:=\frac{\left(R^{h}\right)^{T} \nabla_{h} \tilde{y}^{h}-\mathrm{Id}}{h^{\beta / 2}}
$$

Then up to a subsequence $G^{h} \rightharpoonup G$ weakly in $L^{2}\left(\Omega^{1}, \mathbb{R}^{3 \times 3}\right)$ where

$$
\begin{equation*}
\overparen{G}=G_{0}\left(x^{\prime}\right)+\left(x_{3}-\frac{1}{2}\right) G_{1}\left(x^{\prime}\right) \tag{2.40}
\end{equation*}
$$

with

$$
G_{1}=-\nabla^{2} u
$$

and
(i) if $\beta=4$, then $\operatorname{sym} G_{0}=\operatorname{sym} \nabla w+\frac{1}{2} \nabla u \otimes \nabla u$,
(ii) if $\beta>4$, then $\operatorname{sym} G_{0}=\operatorname{sym} \nabla w$

Proof. The proof is as in [15, Lemma 2]. We include it for the convenience of the reader. We define the difference quotient

$$
H^{h}\left(x^{\prime}, x_{3}, s\right):=\frac{G^{h}\left(x^{\prime}, x_{3}+s\right)-G^{h}\left(x^{\prime}, x_{3}\right)}{s}
$$

Letting

$$
\eta^{h}\left(x^{\prime}, x_{3}, s\right):=\frac{1}{h^{\beta / 2-1}}\left(\frac{1}{s} \int_{0}^{s} \frac{1}{h} \partial_{3}\left(\tilde{y}^{h}\right)^{\prime}\left(x^{\prime}, x_{3}+\sigma\right) d \sigma\right)
$$

We first observe that for all $x_{3}, x_{3}+s \in[0,1]$

$$
\begin{aligned}
\| \eta^{h}\left(x^{\prime}, x_{3}, s\right) & -\left(A_{13}, A_{23}\right)^{T}\left\|_{L^{2}(\Omega)}=\right\| \frac{1}{h^{\beta / 2-1}}\left(\frac{1}{s} \int_{0}^{s} \partial_{3}\left(y^{h}\right)^{\prime}\left(x^{\prime}, h\left(x_{3}+\sigma\right)\right) d \sigma\right)-\left(A_{13}, A_{23}\right)^{T} \|_{L^{2}(\Omega)} \\
& \leq\left\|\frac{1}{h^{\beta / 2-1}}\left(\frac{1}{s} \int_{0}^{s} \partial_{3}\left(y^{h}\right)^{\prime}\left(x^{\prime}, h\left(x_{3}+\sigma\right)\right) d \sigma\right)-\left(A_{13}^{h}, A_{23}^{h}\right)^{T}\right\|_{L^{2}(\Omega)}+\left\|A^{h}-A\right\|_{L^{2}(\Omega)} \\
& \leq\left\|\frac{1}{s} \int_{0}^{s} \frac{1}{h^{\beta / 2-1}}\left(\nabla y^{h}-\mathrm{Id}\right)\left(x^{\prime}, h\left(x_{3}+\sigma\right)\right)-A^{h}\left(x^{\prime}\right) d \sigma\right\|_{L^{2}(\Omega)}+\left\|A^{h}-A\right\|_{L^{2}(\Omega)} \\
& \leq \frac{1}{\sqrt{h s}}\left\|\frac{1}{h^{\beta / 2-1}}\left(\nabla y^{h}-\mathrm{Id}\right)-A^{h}\right\|_{L^{2}\left(\Omega \times\left[h x_{3}, h\left(x_{3}+s\right)\right]\right)}+\left\|A^{h}-A\right\|_{L^{2}(\Omega)} \\
& \leq \frac{1}{\sqrt{h s}}\left\|\frac{1}{h^{\beta / 2-1}}\left(\nabla y^{h}-R^{h}\right)\right\|_{L^{2}\left(\Omega^{h}\right)}+\left\|A^{h}-A\right\|_{L^{2}(\Omega)} \xrightarrow{h \rightarrow 0} 0
\end{aligned}
$$

where we used (2.30) and the definition of $A^{h}$ in (2.31). As a consequence, and in view of (2.38),

$$
\eta^{h}\left(\cdot, x_{3}, s\right) \longrightarrow-\nabla u \quad \text { in } L^{2}(\Omega)
$$

which implies

$$
\nabla^{\prime} \eta^{h}\left(\cdot, x_{3}, s\right) \longrightarrow-\nabla^{2} u \quad \text { in } \mathcal{D}^{\prime}(\Omega)
$$

On the other hand we have by straightforward calculation of the integral in the definition of $\eta^{h}$

$$
\eta^{h}\left(\cdot, x_{3}, s\right)=h^{-\beta / 2} \frac{\left(\tilde{y}^{h}\right)^{\prime}\left(\cdot, x_{3}+s\right)-\left(\tilde{y}^{h}\right)^{\prime}\left(\cdot, x_{3}\right)}{s}
$$

which yields

$$
\nabla^{\prime} \eta^{h}\left(\cdot, x_{3}, s\right)=\overparen{R^{h} H^{h}}\left(x^{\prime}, x_{3}, s\right) \rightharpoonup \frac{\overparen{G}\left(x^{\prime}, x_{3}+s\right)-\overparen{G}\left(x^{\prime}, x_{3}\right)}{s} \quad \text { weakly in } L^{2}(\Omega)
$$

since $R^{h}$ is uniformly bounded and converges to Id in $L^{2}(\Omega)$. Comparing the two expressions of the limit for $\nabla^{\prime} \eta^{h}$ we obtain

$$
\frac{\widehat{G}\left(\cdot, x_{3}+s\right)-\overparen{G}\left(\cdot, x_{3}\right)}{s}=-\nabla^{2} u
$$

Since $\nabla^{2} u$ is independent of $x_{3}$, we obtain that $\overparen{G}$ is affine in $x_{3}$, and can be written in the form given in (2.40), and that $G_{1}=-\nabla^{2} u$. In order to determine sym $G_{0}$, we let

$$
G_{0}^{h}\left(x^{\prime}\right):=\int_{0}^{1} G^{h}\left(x^{\prime}, x_{3}\right) d x_{3}
$$

Note that

$$
G^{h}=\underbrace{\frac{1}{h^{\beta / 2}}\left(\nabla_{h} \tilde{y}^{h}-\mathrm{Id}\right)-\frac{1}{h^{\beta / 2}}\left(R^{h}-\mathrm{Id}\right)}_{\widetilde{G}^{h}}+\underbrace{\frac{1}{h^{\beta / 2}}\left(R^{h}-\mathrm{Id}\right)^{T}\left(\nabla_{h} \tilde{y}^{h}-R^{h}\right)}_{\bar{G}^{h}} .
$$

The last term $\bar{G}^{h}$ converges to 0 in $L^{2}\left(\Omega^{1}\right)$ in view of (2.30). Considering the upper left submatrix of the first two terms we have

$$
\begin{aligned}
\int_{0}^{1} \operatorname{sym} \overparen{\widetilde{G}^{h}}\left(\cdot, x_{3}\right) d x_{3} & \left.=\frac{1}{h^{\beta / 2}} \int_{0}^{1} \operatorname{sym}\left(\nabla^{\prime}\left(\tilde{y}^{h}\right)^{\prime}-\mathrm{Id}\right)\left(\cdot, x_{3}\right) d x_{3}-\frac{1}{h^{\beta / 2}} \operatorname{sym} \overparen{\left(R^{h}\right.}-\mathrm{Id}\right) \\
& =\frac{1}{h^{\beta / 2}} \operatorname{sym} \nabla\left(\frac{1}{h} \int_{0}^{h}\left(\left(y^{h}\right)^{\prime}-\mathrm{id}\right)\left(\cdot, x_{3}\right) d x_{3}\right)-\frac{1}{h^{\beta / 2}} \operatorname{sym}\left(\overparen{R^{h}}-\mathrm{Id}\right) \\
& =\frac{h^{\min \{\beta-2, \beta / 2\}}}{h^{\beta / 2}} \operatorname{sym} \nabla w_{h}-\frac{h^{\beta-2}}{h^{\beta / 2}} \overparen{B^{h}}
\end{aligned}
$$

by (2.29) and (2.34). As a consequence, and in view of Theorem 2.10(iii), (2.34) and (2.39)

$$
\operatorname{sym} \overparen{G_{0}^{h}} \xrightarrow{\text { weakly in } L^{2}(\Omega)} \begin{cases}\operatorname{sym} \nabla w+\frac{1}{2} \nabla u \otimes \nabla u & \text { if } \beta=4 \\ \operatorname{sym} \nabla w & \text { if } \beta>4\end{cases}
$$

But on the other hand the weak limit of $\operatorname{sym} \overparen{G_{0}^{h}}$ must be equal to $\int_{0}^{1} \overparen{G}\left(\cdot, x_{3}\right) d x_{3}$, which is equal to $G_{0}$ in view of (2.40). The proof is complete.

## 2.5. $\Gamma$-liminf estimates.

Theorem 2.13. Let $\Omega, \Omega^{h}, K$ be as above, and assume that $\operatorname{Cap}_{1,2}(K)>0$ and $\beta>2$. Assume that for a sequence as $h \rightarrow 0, y^{h} \in \mathcal{A}_{K}^{h}$,
(H) $u_{h} \rightarrow u$, in $W^{1,2}(\Omega)$, and $w_{h} \rightharpoonup w$ weakly in $W^{1,2}(\Omega)$,
where $u_{h}$ and $w_{h}$ are respectively defined as in (2.28) and (2.29). Then
(i) If $2<\beta<4$

$$
\underset{h}{\liminf } \frac{1}{h^{\beta}} J^{h}\left(y^{h}\right) \geq \int_{\Omega} \frac{1}{24} Q_{2}\left(\nabla^{2} u\right)-u f d x .
$$

(ii) If $\beta=4$

$$
\underset{h}{\liminf } \frac{1}{h^{\beta}} J^{h}\left(y^{h}\right) \geq \int_{\Omega} \frac{1}{2} Q_{2}\left(\operatorname{sym} \nabla w+\frac{1}{2} \nabla u \otimes \nabla u\right)+\frac{1}{24} Q_{2}\left(\nabla^{2} u\right)-u f d x .
$$

(iii) If $\beta>4$

$$
\liminf _{h} \frac{1}{h^{\beta}} J^{h}\left(y^{h}\right) \geq \int_{\Omega} \frac{1}{2} Q_{2}(\operatorname{sym} \nabla w)+\frac{1}{24} Q_{2}\left(\nabla^{2} u\right)-u f d x .
$$

Proof. We first note that under the hypothesis (H) and the scaling the linear term of the energies converge, i.e.

$$
-\frac{1}{h} \int_{\Omega^{h}} \mathbf{f}^{h} \cdot y^{h} d x \xrightarrow{h \rightarrow 0} \int_{\Omega}-u f d x
$$

Therefore, in case $\liminf _{h} \frac{1}{h^{\beta}} E^{h}\left(y^{h}\right)=+\infty$, the results follow trivially. Otherwise, for a suitable subsequence and a constant $C>0$, we obtain $E^{h}\left(y^{h}\right) \leq C h^{\beta}$, and hence the hypotheses of Theorem 2.10 are satisfied. Note that the limits $u, w$ obtained in Theorem 2.10 must be the same given under the convergence hypothesis (H), and so we obtain that $u \in H_{0, K}^{2}(\Omega)$. Applying

Corollary 2.12 , the rest of the proof follows exactly as in [15, Corollary 2$]$, and is left to the reader.
2.6. Recovery sequence and $\Gamma$-limsup estimate for $\beta>4$. In this section we explain how to construct a recovery sequence. Namely, we want to prove the following statement.

Theorem 2.14. Let $\beta>4$. For every $u \in H_{0, K}^{2}(\Omega), w \in H_{0, K}^{1}(\Omega)$, there exists a sequence $y^{h} \in \mathcal{A}_{K}^{h}$ such that for the sequences $u_{h}, w_{h}$ defined as in (2.28) and (2.29), $\left(u_{h}, w_{h}\right) \longrightarrow(u, w)$ in $W^{1,2}(\Omega)$, and

$$
\limsup _{h} \frac{1}{h^{\beta}} J^{h}\left(y^{h}\right) \leq \int_{\Omega} \frac{1}{2} Q_{2}(\operatorname{sym} \nabla w)+\frac{1}{24} Q_{2}\left(\nabla^{2} u\right)-u f d x
$$

Proof. The proof starts from the construction given in [15, Section 6.2]. More precisely we take $\beta=2 \alpha-2$ fix given $u \in H_{0, K}^{2}(\Omega), w \in H_{0, K}^{1}(\Omega)$, and $g \in L^{2}$ a given vector field. Then we define

$$
y^{h}\left(x^{\prime}, x_{3}\right)=\binom{x^{\prime}}{h\left(x_{3}-\frac{1}{2}\right)}+\binom{h^{\alpha-1} w}{h^{\alpha-2} u}-h^{\alpha-1}\left(x_{3}-\frac{1}{2}\right)\left(\begin{array}{c}
\partial_{1} u \\
\partial_{2} u \\
0
\end{array}\right)+\frac{h^{\alpha}}{2}\left(x_{3}-\frac{1}{2}\right)^{2} g
$$

The required convergence of $\left(u_{h}, w_{h}\right)$ is established by a straightforward computation. Also, computations as in [15] imply that

$$
\frac{1}{h^{\beta}} W\left(\nabla y^{h}\right) \xrightarrow{h \rightarrow 0} \frac{1}{2} Q_{3}\left(\operatorname{sym} \nabla w+x_{3} B\right)=\frac{1}{2} Q_{3}(\operatorname{sym} \nabla w)+\frac{1}{2} x_{3}(\operatorname{sym} \nabla w: B)+\frac{1}{2} Q_{3}\left(x_{3} B\right)
$$

where

$$
B:=-\nabla^{2} u+\operatorname{sym}\left(g \otimes e_{3}\right)
$$

As observed in [15], one can choose $g$ in such a way that

$$
\int_{\Omega \times\left(-\frac{1}{2}, \frac{1}{2}\right)} Q_{3}\left(x_{3} B\right)=\frac{1}{12} \int_{\Omega} Q_{2}\left(\nabla^{2} u\right)
$$

This choice yields

$$
\frac{1}{h^{\beta}} \int_{\Omega \times\left(-\frac{1}{2}, \frac{1}{2}\right)} W\left(\nabla y^{h}\right) \xrightarrow{h \rightarrow 0} \frac{1}{2} \int_{\Omega} Q_{3}(\operatorname{sym} \nabla w)+\frac{1}{24} \int_{\Omega} Q_{2}\left(\nabla^{2} u\right)
$$

Now we see in the direct expression of $y^{h}$ and the fact that $u, w \in H_{0, K}^{1}(\Omega)$ implies that $y^{h}=0$ on $K$. This is not enough to have $y^{h} \in \mathcal{A}_{K}^{h}$ because we need $y^{h}=0$ on $K_{h}$, a neighborhood of $K$. Therefore we need first to approximate $u, w$ by some $u_{h}, w_{h}$ which are equal to 0 on $K_{h}$. This can easily be done by use of a sort of diagonal argument. Indeed, $u \in H_{0, K}^{2}(\Omega), w \in H_{0, K}^{1}(\Omega)$, and thus from Lemma A. 1 we know that there exist sequences $u_{n}, w_{n} \in C^{\infty}(\bar{\Omega})$ such that $\operatorname{supp}\left(u_{n}\right) \cup \operatorname{supp}\left(w_{n}\right) \cap K=\emptyset, u_{n} \rightarrow u$ in $W^{2,2}(\Omega)$, and $w_{n} \rightarrow w$ in $W^{1,2}(\Omega)$. Since for $n$ fixed, $K_{h} \subset \Omega \backslash\left(\operatorname{supp}\left(u_{n}\right) \cup \operatorname{supp}\left(w_{n}\right)\right)$ for $h$ small, we can then define a subsequence $n_{h} \rightarrow+\infty$ such that $K_{h} \subset \Omega \backslash\left(\operatorname{supp}\left(u_{n_{h}}\right) \cup \operatorname{supp}\left(w_{n_{h}}\right)\right)$. By this way if we change $y^{h}$ as above with $u_{n_{h}}$, $w_{n_{h}}$ instead of $u, w$, then we have now $y^{h} \in \mathcal{A}_{K}^{h}$. Since the convergences of $u_{n_{h}}$ to $u$, and $w_{n_{h}}$ to $w$, hold respectively in the strong topologies of $W^{2,2}$ and $W^{1,2}$, it is easy to verify that we still have the same limit

$$
\frac{1}{h^{\beta}} W\left(\nabla y^{h}\right) \xrightarrow{h \rightarrow 0} \frac{1}{2} Q_{3}\left(\operatorname{sym} \nabla w+x_{3} B\right),
$$

which finishes the proof of the proposition.

Remark 2.15. Theorems 2.13 and 2.14 mean that for $\beta>4$, the functionals $h^{-\beta} E^{h}, \Gamma$-converge to the functional

$$
I(u, w):=\int_{\Omega} \frac{1}{2} Q_{2}(\operatorname{sym} \nabla w)+\frac{1}{24} Q_{2}\left(\nabla^{2} u\right)-u f d x
$$

defined over $H_{0, K}^{2} \times H_{0, K}^{1}$, under the strong convergence of $\left(u_{h}, w_{h}\right)$ to $(u, w)$ in $W^{2,2} \times W^{1,2}$. Note that the strong convergence of $w_{h}$, and hence the inclusion of the weak limit $w$ in $H_{0, K}^{1}$ in Theorem 2.10 is not established in its full generality for this regime. To bypass this problem, and in view of the fact that $u$ and $w$ are decoupled in the expression of $I(u, w)$, one can forgo the mention of $w$ in the final result. Indeed, notice that both inequalities still work for $w=0$, i.e.

$$
\liminf _{h} \frac{1}{h^{\beta}} J^{h}\left(y^{h}\right) \geq \frac{1}{24} Q_{2}\left(\nabla^{2} u\right)-u f d x
$$

and the same recovery sequence written for $u \in H_{0, K}^{2}(\Omega)$ and $w=0$ satisfies

$$
\limsup _{h} \frac{1}{h^{\beta}} J^{h}\left(y^{h}\right) \leq \frac{1}{24} Q_{2}\left(\nabla^{2} u\right)-u f d x
$$

As stated in Theorem 1.1, this establishes the functional

$$
I(u):=\int_{\Omega} \frac{1}{24} Q_{2}\left(\nabla^{2} u\right)-u f d x
$$

over $H_{0, K}^{2}$ as the $\Gamma$-limit of $h^{-\beta} E^{h}$ under the strong convergence of $u_{h}$ to $u$.
2.7. Energy scalings and $\beta$-minimizing sequences. We first prove the following proposition regarding the fact that the scaling of the infimum energy is determined by the scaling of body forces in our setting. Remember that vertical body forces $\mathbf{f}^{h}: \Omega^{h} \rightarrow \mathbb{R}^{3}$ are assumed to be of the form $\mathbf{f}^{h}:=\left(0,0, h^{\alpha} \tilde{f}\right)$, for $\tilde{f} \in L^{2}(\Omega)$.

Proposition 2.16. Assume $\alpha>2$, and let $\beta=2 \alpha-2$. There exists $h_{0}>0$ such that for $a$ constant $C:=C\left(\Omega, K, \tilde{f}, h_{0}\right)>0$ and all $h<h_{0}$

$$
-C h^{\beta} \leq \inf _{\mathcal{A}_{K}^{h}} J^{h} \leq C h^{\alpha} \ll C h^{\beta}
$$

In particular, since $\inf _{\mathcal{A}_{K}^{h}} J^{h} \in \mathbb{R}$, there exists a $\beta$-minimizing sequence.
Proof. First note that letting $y=\mathrm{id}_{\Omega^{h}}$ the identity map we have

$$
J^{h}\left(\operatorname{id}_{\Omega^{h}}\right)=-\frac{1}{h} \int_{\Omega^{h}} \mathbf{f}^{h}(x) \cdot x d x \leq C h^{\alpha}\|\tilde{f}\|_{L^{2}(\Omega)}
$$

yielding the upper bound on $\inf _{\mathcal{A}_{K}^{h}} J^{h}$.
Let now $y \in \mathcal{A}_{K}^{h}$ be an arbitrary deformation, and consider the matrix field $\widetilde{F}$ obtained from Lemma 2.6. Letting $\widetilde{F}_{0}:=f_{\Omega} \widetilde{F} d x$. We have by the Poincaré inequality

$$
\left\|\widetilde{F}-\widetilde{F}_{0}\right\|_{L^{2}(\Omega)}^{2} \leq C\|\nabla \widetilde{F}\|_{L^{2}(\Omega)}^{2} \leq \frac{C}{h^{2}} E^{h}(y)
$$

which implies

$$
\frac{1}{h}\left\|\nabla y-\widetilde{F}_{0}\right\|_{L^{2}\left(\Omega^{h}\right)}^{2} \leq 2\left(\frac{1}{h}\|\nabla y-\widetilde{F}\|_{L^{2}\left(\Omega^{h}\right)}^{2}+\frac{1}{h}\left\|F-\widetilde{F}_{0}\right\|_{L^{2}\left(\Omega^{h}\right)}^{2}\right) \leq \frac{C}{h^{2}} E^{h}(y)
$$

Note that

$$
\left|\widetilde{F}_{0}\right|^{2} \leq f_{\Omega}|\widetilde{F}|^{2} d x \leq C\left(f_{\Omega} \operatorname{dist}^{2}(\tilde{F}, S O(3)) d x+1\right) \leq C\left(\|\tilde{d}\|_{L^{2}(\Omega)}+1\right) \leq C\left(E^{h}(y)+1\right)
$$

Combining the last two estimates we obtain

$$
\frac{1}{h}\|\nabla y\|_{L^{2}\left(\Omega^{h}\right)}^{2} \leq C\left(\frac{1}{h^{2}} E^{h}(y)+1\right)
$$

Now, since $y \in \mathcal{A}_{K}^{h}$, we have $y_{3}=0$ on $K_{h}$, i.e $y_{3} \in \mathbf{A}_{K}^{h}$ and we can apply Theorem A. 5 to obtain

$$
\left\|y_{3}\right\|_{L^{2}\left(\Omega^{h}\right)}^{2} \leq C\left\|\nabla y_{3}\right\|_{L^{2}\left(\Omega^{h}\right)}^{2}
$$

Therefore

$$
\frac{1}{h}\left\|y_{3}\right\|_{L^{2}\left(\Omega^{h}\right)}^{2} \leq \frac{C}{h}\left\|\nabla y_{3}\right\|_{L^{2}\left(\Omega^{h}\right)}^{2} \leq C\left(\frac{1}{h^{2}} E^{h}(y)+1\right)
$$

and hence

$$
\begin{aligned}
J^{h}(y) & =E^{h}(y)-\frac{1}{h} \int_{\Omega^{h}} \mathbf{f}^{h} \cdot y d x \geq E^{h}(y)-C h^{\alpha}\|\tilde{f}\|_{L^{2}(\Omega)}\left(\frac{1}{\sqrt{h}}\left\|y_{3}\right\|_{L^{2}\left(\Omega^{h}\right)}\right) \\
& \geq E^{h}(y)-C h^{\alpha}\|\tilde{f}\|_{L^{2}(\Omega)}\left(\frac{1}{h}\left(E^{h}(y)\right)^{\frac{1}{2}}+1\right) \\
& =E^{h}(y)-C h^{\alpha-1}\|\tilde{f}\|_{L^{2}(\Omega)} E^{h}(y)^{\frac{1}{2}}-C h^{\alpha}\|\tilde{f}\|_{L^{2}(\Omega)} .
\end{aligned}
$$

Hence

$$
\begin{equation*}
J^{h}(y) \geq\left(E^{h}(y)^{\frac{1}{2}}-C h^{\alpha-1}\|\tilde{f}\|_{L^{2}(\Omega)}\right)^{2}-C^{2} h^{2 \alpha-2}\|\tilde{f}\|_{L^{2}(\Omega)}^{2}-C h^{\alpha}\|\tilde{f}\|_{L^{2}(\Omega)} \tag{2.41}
\end{equation*}
$$

A standard argument now yields for an adjusted constant $C$

$$
J^{h}(y) \geq-C h^{2 \alpha-2}\|\tilde{f}\|_{L^{2}(\Omega)} \geq-C h^{\beta}
$$

implying the required lower bound on $\inf _{\mathcal{A}_{K}^{h}} J^{h}$.
Corollary 2.17. Let $\alpha>2$ and $\beta=2 \alpha-2$. For any sequence $y^{h} \in \mathcal{A}_{K}^{h}$,

$$
J^{h}\left(y^{h}\right) \precsim h^{\beta} \Longrightarrow E^{h}\left(y^{h}\right) \precsim h^{\beta} .
$$

In particular any $\beta$-minimizing sequence $y^{h}$ satisfies $E^{h}\left(y^{h}\right) \leq C h^{\beta}$.
Proof. Assuming $J^{h}\left(y^{h}\right) \precsim h^{\beta}$, from (2.41) we obtain

$$
\left(E^{h}\left(y^{h}\right)^{\frac{1}{2}}-C h^{\beta / 2}\right)^{2} \leq J^{h}\left(y^{h}\right)+C h^{\beta}+C h^{\alpha} \leq C h^{\beta}
$$

which implies $E^{h}\left(y^{h}\right) \precsim h^{\beta}$. Consider now a $\beta$-minimizing sequence $y^{h}$,

$$
\limsup _{h \rightarrow 0} \frac{1}{h^{\beta}}\left(J^{h}\left(y^{h}\right)-\inf _{\mathcal{A}_{K}^{h}} J^{h}\right)=0
$$

We have therefore for $h$ small enough

$$
J^{h}\left(y^{h}\right) \leq\left(\inf _{\mathcal{A}_{K}^{h}} J^{h}+1\right) h^{\beta} \leq C h^{\beta}
$$

which implies the required estimate on $E^{h}\left(y^{h}\right)$.
Theorem 2.18. Let $\alpha>3, \beta=2 \alpha-2>4$, and let $y^{h} \in \mathcal{A}_{K}^{h}$ be a $\beta$-minimizing sequence for $J^{h}$, and let $\left(u_{h}, w_{h}\right)$ be defined as in (2.28), (2.29). Then up to a subsequence, $\left(u_{h}, w_{h}\right) \longrightarrow(u, 0)$ in $W^{1,2}(\Omega), u \in H_{0, K}^{2}(\Omega)$, and $u$ minimizes the functional

$$
I(u):=\int_{\Omega} \frac{1}{24} Q_{2}\left(\nabla^{2} u\right)-u f d x
$$

over $H_{0, K}^{2}(\Omega)$. Conversely, if $u$ minimizes $I(u)$ over $H_{0, K}^{2}(\Omega)$, then there exists a $\beta$-minimizing sequence $y^{h} \in \mathcal{A}_{K}^{h}$ for $J^{h}$ such that $\left(u_{h}, w_{h}\right) \longrightarrow(u, 0)$ in $W^{1,2}(\Omega)$.

Proof. Combining Corollary 2.17 and Theorem 2.10, we obtain that up to a subsequence $u_{h} \longrightarrow u$ in $W^{1,2}(\Omega), u \in H_{0, K}^{2}(\Omega)$, and that $w_{h} \rightharpoonup w$ weakly in $W^{1,2}\left(\Omega, \mathbb{R}^{2}\right)$. A standard argument using the $\Gamma$-convergence results Theorem 2.13 and Theorem 2.14 (see also Remark 2.15) shows that $u$ minimize the functional $I(u)$ over $H_{0, K}^{2}(\Omega)$, and that

$$
I(u)=\min _{u \in H_{0, K}^{2}(\Omega)} I=\lim _{h \rightarrow 0} \frac{1}{h^{\beta}} J^{h}\left(y^{h}\right)=\lim _{h \rightarrow 0} \frac{1}{h^{\beta}} \inf _{\mathcal{A}_{K}^{h}} J^{h}
$$

On the other hand, Theorem 2.13 implies that the recovery sequence $y^{h}$ of Theorem 2.14 for $w=0$, which satisfies the required convergence criteria, satisfies

$$
\lim _{h \rightarrow 0} \frac{1}{h^{\beta}} J^{h}\left(y^{h}\right)=I(u)=\lim _{h \rightarrow 0} \frac{1}{h^{\beta}} \inf _{\mathcal{A}_{K}^{h}} J^{h},
$$

which implies that $y^{h}$ is $\beta$-minimizing.
In fact, more can be shown with a finer reasoning. Following the same arguments as in [15, Section 7.2], the equi-integrability of the family $\mathcal{Y}$ and the strong convergence of $\operatorname{sym} \nabla w_{h}$ in $L^{2}$ can be established. As a consequence the strong convergence of $w_{h}$ to $w \in H_{0, K}^{1}$ follows immediately from Theorem 2.10-(v). Using this stronger convergence, and the full force of Theorems 2.13 and 2.14, it can be shown that the pair $(u, w)$ must minimize the functional

$$
\int_{\Omega} \frac{1}{2} Q_{2}(\operatorname{sym} \nabla w)+I(u)
$$

over $H_{0, K}^{2} \times H_{0, K}^{1}$. Since $u$ and $w$ are decoupled in the energy, we must have sym $\nabla w=0$ for the minimizer $(u, w)$, and hence the infinitesimal rigidity of displacements in $\mathbb{R}^{2}$ implies that $w(x)=D x+b$ is an affine map with $D \in s o(2)$, skew-symmetric, and therefore $\operatorname{det} D \neq 0$ or $D=0$. Since $w \in H_{0, K}^{1}$ and $\operatorname{Cap}_{1,2}(K)>0$, we deduce that $w=0$, since otherwise $w$ can vanish at only one single point. We conclude that $w_{h} \rightarrow 0$ in $W^{1,2}(\Omega)$.

## 3. The optimal biharmonic support problem

In this section we know focus on shape optimisation Problem in (1.6) and we assume for simplicity that $\partial \Omega \subset K$. By this way the space $H_{0, K}^{k}(\Omega)$ reduces to the more classical $H_{0}^{k}(\Omega \backslash K)$. Our aim is to prove existence and regularity of minimizers $K$ for that problem.
3.1. Dual formulation. Let $\Omega \subset \mathbb{R}^{2}$ be any open and simply connected domain. We denote by $\mathcal{K}(\Omega)$ the collection of all compact connected subsets $K \subset \bar{\Omega}$ such that $\partial \Omega \subset K$. In particular, for all $K \in \mathcal{K}(\Omega)$, any connected component of $\Omega \backslash K$ is necessarily simply connected.

Notice that for $u \in C_{c}^{\infty}(\Omega)$ it is very classical that

$$
\int_{\Omega}\left|\nabla^{2} u\right|^{2} d x=\int_{\Omega}|\Delta u|^{2} d x
$$

Indeed, a simple computation yields

$$
\begin{align*}
\int_{\Omega}\left|\Delta^{2} u\right|^{2} d x & =\int_{\Omega}|\Delta u|^{2} d x+2 \int_{\Omega} u_{12}^{2}-u_{11} u_{22} d x \\
& =\int_{\Omega}|\Delta u|^{2} d x+2 \int_{\Omega} \operatorname{div} F d x \tag{3.1}
\end{align*}
$$

where

$$
F:=\binom{u_{12} u_{2}-u_{22} u_{1}}{u_{12} u_{1}-u_{11} u_{2}}
$$

Now

$$
\int_{\Omega} \operatorname{div} F d x=0
$$

because $F \in C_{c}^{\infty}(\Omega)$. Therefore, the biharmonic equation can be equivalently solved either by minimizing $\int_{\Omega \backslash K}\left|\nabla^{2} u\right|^{2} d x$ or $\int_{\Omega \backslash K}|\Delta u|^{2} d x$ in $H_{0}^{2}(\Omega \backslash K)$. We denote by $L_{s y m}^{2}(\Omega)$ the space of symmetric matrix valued $L^{2}$ functions that we endow with the scalar product

$$
\langle A, M\rangle=\int_{\Omega} A: M d x
$$

We denote by $\operatorname{div} M$ the divergence of the Matrix valued function $M$, which consists in taking the divergence of each raw of $M$. A simple computation reveals that for $u \in C_{c}^{\infty}(\Omega)$ and $M \in C_{\text {sym }}^{\infty}(\Omega)$ it holds

$$
\int_{\Omega} \nabla^{2} u: M d x=\int_{\Omega} u \operatorname{divdiv} M d x
$$

which naturally extends to $u \in H_{0}^{2}(\Omega)$ and $M \in H_{\text {sym }}^{2}(\Omega)$.
We begin with an elementary Lemma.
Lemma 3.1. Let $f: E \times F \rightarrow \mathbb{R}$ be a real valued function defined on two given sets $E, F$. Assume that $\left(u_{0}, v_{0}\right)$ is a saddle point, i.e. satisfies

$$
f\left(u_{0}, v_{0}\right)=\max _{v} f\left(u_{0}, v\right)=\min _{u} f\left(u, v_{0}\right)
$$

Then

$$
\inf _{u} \sup _{v} f(u, v)=\sup _{v} \inf _{u} f(u, v)=f\left(u_{0}, v_{0}\right)
$$

Proof. For any $(u, w)$ fixed it is clear that $f(u, w) \leq \sup _{v} f(u, v)$ thus taking first the inf in $u$ and then sup in $w$ yields

$$
\sup _{w} \inf _{u} f(u, w) \leq \inf _{u} \sup _{v} f(u, v)
$$

For the reverse inequality we use that $\left(u_{0}, v_{0}\right)$ is a saddle point. Indeed,

$$
\inf _{u} \sup _{v} f(u, v) \leq \max _{v} f\left(u_{0}, v\right)=\min _{u} f\left(u, v_{0}\right) \leq \sup _{v} \inf _{u} f(u, v)
$$

which concludes the proof.
Proposition 3.2. Let $F: H_{0}^{2}(\Omega \backslash K) \times L_{\text {sym }}^{2}(\Omega) \rightarrow \mathbb{R}$ be defined by

$$
F(u, M):=2 \int_{\Omega} \nabla^{2} u: M d x-\int_{\Omega}|M|^{2} d x-2 \int_{\Omega} u f d x
$$

Then $\left(u_{K}, \nabla^{2} u_{K}\right)$ is a sadle point for $F$. In particular,

$$
-\int_{\Omega}\left|\nabla^{2} u_{K}\right|^{2} d x=\sup _{M \in L_{s y m}^{2}(\Omega)} \inf _{u \in H_{0}^{2}(\Omega \backslash K)} F(u, M)
$$

Proof. From the elementary equality $|A+B|^{2}=|A|^{2}+|B|^{2}+2 A: B$ valid for any matrix $A, B$ we infer that

$$
|A|^{2} \geq 2 A: B-|B|^{2}
$$

and in particular for any $M \in L_{\text {sym }}^{2}(\Omega)$, we always have

$$
F\left(u_{K}, \nabla^{2} u_{K}\right) \geq F\left(u_{K}, M\right)
$$

and since the equality occurs for $M=\nabla^{2} u_{K}$ we can affirm

$$
\begin{equation*}
F\left(u_{K}, \nabla^{2} u_{K}\right)=\max _{M} F\left(u_{K}, M\right) . \tag{3.2}
\end{equation*}
$$

On the other hand, by Lax-Milgram theory we know that $u_{K}$ is the unique minimizer in $H_{0}^{2}(\Omega \backslash K)$ for the functional $J$ defined by

$$
J(u):=\int_{\Omega}\left|\nabla^{2} u\right| d x-2 \int_{\Omega} u f d x,
$$

and the weak formulation for this problem says that

$$
\int_{\Omega} \nabla^{2} u: \nabla^{2} u_{K} d x=\int_{\Omega} u f d x \quad \forall u \in H_{0}^{2}(\Omega \backslash K) .
$$

In turn, for $\nabla^{2} u_{K}$ fixed and an arbitrary $u \in H_{0}^{2}(\Omega \backslash K)$, we observe that the expression of $F\left(u, \nabla^{2} u_{K}\right)$ reduces to

$$
F\left(u, \nabla^{2} u_{K}\right)=-\int_{\Omega}\left|\nabla^{2} u_{K}\right|^{2} d x
$$

which is constant in the $u$ variable. We deduce a fortiori that

$$
F\left(u_{K}, \nabla^{2} u_{K}\right)=\min _{u} F\left(u, \nabla^{2} u_{K}\right),
$$

which together with (3.2), proves that $\left(u_{K}, \nabla^{2} u_{K}\right)$ is a saddle point for $F(u, M)$. We conclude by applying Lemma 3.1.

In the following proposition we will use the notation $L D(A)$ for a function $v \in L^{2}\left(A, \mathbb{R}^{2}\right)$ satisfying $e(v) \in L^{2}(A)$, where $e(v)=\left(\nabla v+\nabla v^{T}\right) / 2$ is the symmetrized gradient of $v$.

Proposition 3.3. Let $\varphi \in H_{0}^{1}(\Omega)$ be the unique solution of

$$
\left\{\begin{array}{c}
-\Delta \varphi=f \text { in } \Omega \\
\varphi \in H_{0}^{1}(\Omega)
\end{array}\right.
$$

and let

$$
G=\varphi \operatorname{Id}=\left(\begin{array}{ll}
\varphi & 0 \\
0 & \varphi
\end{array}\right)
$$

Let $K$ be a minimizer for Problem (1.6). Then there exists $v_{K} \in L D(B(x, r) \backslash K)$ such that $\left(v_{K}, K\right)$ is a minimizer for the problem

$$
\begin{equation*}
\min _{(v, K) \in \mathcal{A}} \int_{\Omega \backslash K}|e(v)-G|^{2} d x+\mathcal{H}^{1}(K) \tag{3.3}
\end{equation*}
$$

where

$$
\mathcal{A}:=\{(v, K): K \in \mathcal{K}(\Omega) \text { and } v \in L D(B(x, r) \backslash K)\} .
$$

Proof. For a given $K \subset \bar{\Omega}$, we can apply Proposition 3.2 to write

$$
\begin{aligned}
-\int_{\Omega}\left|\nabla^{2} u_{K}\right|^{2} d x & =\sup _{M \in L_{s y m}^{2} u \in H_{0}^{2}(\Omega \backslash K)} \inf _{\Omega} 2 \int_{\Omega} \nabla^{2} u: M d x-\int_{\Omega}|M|^{2} d x-2 \int_{\Omega} u f d x \\
& =\sup _{M \in L_{s y m}^{2} u \in H_{0}^{2}(\Omega \backslash K)} \inf _{H_{0}^{2}}\langle u, \operatorname{div} \operatorname{div} M\rangle_{\left(H_{0}^{2}\right)^{\prime}}-\int_{\Omega}|M|^{2} d x-2 \int_{\Omega} u f d x . \\
& =\sup _{M \in L_{s y m}^{2} u \in H_{0}^{2}(\Omega \backslash K)} \inf _{H_{0}^{2}}\langle u, \operatorname{div} \operatorname{div} M-f\rangle_{\left(H_{0}^{2}\right)^{\prime}}-\int_{\Omega}|M|^{2} d x .
\end{aligned}
$$

The infimum in the $u$ variable in the above is equal to $-\infty$, excepted when $\operatorname{div} \operatorname{div} M=f$ in $\mathcal{D}^{\prime}(\Omega \backslash K)$. This leads to the following dual equality

$$
\int_{\Omega}\left|\nabla^{2} u_{K}\right|^{2} d x=\min \left\{\int_{\Omega}|M|^{2} d x \quad \text { s.t. } M \in L_{\text {sym }}^{2}(\Omega) \text { and } \operatorname{div} \operatorname{div} M=f \text { in } \mathcal{D}^{\prime}(\Omega \backslash K)\right\} .
$$

Moreover, the minimium is achieved for $M=\nabla^{2} u_{K}$. In other words

$$
\min _{K} \int_{\Omega}\left|\nabla^{2} u_{K}\right|^{2} d x+\mathcal{H}^{1}(K)=\min _{K} \min _{M} \int_{\Omega}|M|^{2} d x+\mathcal{H}^{1}(K)
$$

where the minimum in $M$ is over all $M \in L_{\text {sym }}^{2}(\Omega)$ satisfying div $\operatorname{div} M=f$ in $\mathcal{D}^{\prime}(\Omega \backslash K)$.
Now let $v \in H_{0}^{1}(\Omega)$ be the unique solution of

$$
\left\{\begin{array}{c}
-\Delta v=f \text { in } \Omega \\
v \in H_{0}^{1}(\Omega)
\end{array}\right.
$$

and let

$$
G=v \operatorname{Id}=\left(\begin{array}{cc}
v & 0 \\
0 & v
\end{array}\right)
$$

Then $G$ is symetric and $\operatorname{div} \operatorname{div} G=\Delta v=-f$ in $\mathcal{D}^{\prime}(\Omega)$. It follows that for all $M$ satisfying $\operatorname{div} \operatorname{div} M=f$ in $\mathcal{D}^{\prime}(\Omega \backslash K)$,

$$
\operatorname{div} \operatorname{div}(M+G)=0 \text { in } \mathcal{D}^{\prime}(\Omega \backslash K)
$$

In particular, since $\Omega \backslash K$ is simply connected, there exists $u$ such that

$$
\operatorname{div}(M+G)=\nabla^{\perp} u
$$

We deduce that

$$
\operatorname{curl}\left(\operatorname{div}(\mathrm{M}+\mathrm{G})^{\perp}\right)=0
$$

or differently

$$
\operatorname{curlcurl}(\operatorname{Com}(M+G))=0 .
$$

The classical Saint-Venant compatibility condition yields the existence of $u$ such that $\operatorname{Com}(M+$ $G)=e(u)$ in the simply connected domain $\Omega \backslash K$ and we infer that

$$
M=\operatorname{Com}(e(u))-G,
$$

and finally since $\operatorname{Com}(G)=G$,

$$
\int_{\Omega \backslash K}|M|^{2} d x=\int_{\Omega \backslash K}|e(u)-G|^{2} d x .
$$

We conclude that

$$
\min _{M} \int_{\Omega \backslash K}|M|^{2} d x=\min _{u} \int_{\Omega \backslash K}|e(u)-G|^{2} d x
$$

which ends the proof of the proposition.
Remark 3.4. Notice that $G$ is bounded and we actually have

$$
\|G\|_{L^{\infty}(\Omega)} \leq C\|f\|_{p}
$$

Indeed, by elliptic regularity $G \in W^{2, p}(\Omega)$ and then since $p>2$ we know that $u \in C^{1, \alpha}(\bar{\Omega})$.
3.2. Existence. We start by proving the existence of a minimizer, as a simple consequence of Sverak's [23] result.
Proposition 3.5. Let $\Omega$ be an open and bounded set in $\mathbb{R}^{2}$ and let $f \in L^{\infty}(\Omega)$. Let $\left(K_{n}\right)_{n}$ be a sequence of closed connected subset of $\bar{\Omega}$, converging to a closed connected set $K \subset \bar{\Omega}$ with respect to the Hausdorff distance. Then

$$
u_{K_{n}} \xrightarrow{n \rightarrow+\infty} u_{K} \text { strongly in } H^{2}(\Omega) .
$$

Proof. We proceed as a standard $\Gamma$-convergence argument. We start by noticing that for all $n$, using that $u_{K_{n}}$ is the solution of $\Delta^{2} u_{K_{n}}=f$ in $H_{0}^{2}\left(\Omega \backslash K_{n}\right)$,

$$
\int_{\Omega}\left|\nabla^{2} u_{K_{n}}\right|^{2} d x=\int_{\Omega} f u_{K_{n}} \leq\|f\|_{2}^{2}\left\|u_{K_{n}}\right\|_{L^{2}} \leq C\left\|\nabla^{2} u\right\|_{2},
$$

where we have used the Poincaré inequality in $H_{0}^{2}(\Omega)$. This leads to

$$
\int_{\Omega}\left|\nabla^{2} u_{K_{n}}\right|^{2} d x \leq\|f\|_{2}
$$

so that $u_{K_{n}}$ is uniformly bounded in $H_{0}^{2}(\Omega)$. We can therefore extract a subsequence (not relabeled) that converges weakly in $H^{2}(\Omega)$, strongly in $H_{0}^{1}(\Omega)$, and uniformly in $\Omega$ to some function $v \in H_{0}(\Omega)$. By uniform convergence we also know that $v=0$ on the set $K$, the Hausdorff limit of $K_{n}$.

Since $\nabla u_{K_{n}} \in H_{0}^{1}\left(\Omega \backslash K_{n}\right)$ and converges weakly in $H_{0}^{1}(\Omega)$ to $\nabla v$, and since the sets $K_{n}$ are all compact and connected in $\bar{\Omega}$, it follows from [23] that $\nabla v \in H_{0}^{1}(\Omega)$. In other words $v \in H_{0}^{2}(\Omega \backslash K)$.

Now let $\varphi \in C_{c}^{\infty}(\Omega \backslash K)$. By Hausdorff convergence of $K_{n}$ we deduce that the support of $\varphi$ is contained in $\Omega \backslash K_{n}$ for all $n$ large enough. Therefore, we can apply the weak formulation of the problem satisfied by $u_{K_{n}}$ which yields

$$
\int_{\Omega} \nabla^{2} u_{K_{n}}: \nabla^{2} \varphi d x=\int_{\Omega} f \varphi d x .
$$

Passing to the limit we obtain that $u_{K_{n}}$ is the unique solution of $\Delta^{2} u=f$ in $H_{0}^{2}(\Omega \backslash K)$. In other words $u=u_{K}$.

By the week convergence of the Hessians we already have

$$
\int_{\Omega}\left|\nabla^{2} u\right| d x \leq \liminf \int_{\Omega}\left|\nabla^{2} u_{K_{n}}\right|^{2} d x .
$$

Conversely, by definition of $H_{0}^{2}$ there exists a sequence of $v_{n} \in C_{c}^{\infty}(\Omega \backslash K)$ such that $v_{n} \rightarrow u$ strongly in $H^{2}(\Omega)$. Let us first fix a $k_{0} \geq 0$. Using the Hausdorff convergence of $K_{n}$ to $K$ we know that $v_{k_{0}} \in H_{0}^{2}\left(\Omega \backslash K_{n}\right)$ for $n$ large. Since $u_{K_{n}}$ is a minimizer in this class we obtain

$$
\int_{\Omega}\left|\nabla^{2} u_{K_{n}}\right|^{2} d x-2 \int_{\Omega} u_{K_{n}} f d x \leq \int_{\Omega}\left|\nabla^{2} v_{k_{0}}\right|^{2} d x-2 \int_{\Omega} v_{k_{0}} f d x .
$$

passing to the limsup in $n$ we arrive at

$$
\limsup _{n \rightarrow+\infty} \int_{\Omega}\left|\nabla^{2} u_{K_{n}}\right|^{2} d x \leq \int_{\Omega}\left|\nabla^{2} v_{k_{0}}\right|^{2} d x-2 \int_{\Omega} v_{k_{0}} f d x+2 \int_{\Omega} u f d x .
$$

Letting now $k_{0} \rightarrow+\infty$ yields

$$
\limsup _{n \rightarrow+\infty} \int_{\Omega}\left|\nabla^{2} u_{K_{n}}\right|^{2} d x \leq \int_{\Omega}\left|\nabla^{2} u\right|^{2} d x
$$

which together with the liminf inequality proves the strong convergence of $u_{K_{n}}$ in $H^{2}(\Omega)$ to $u=u_{K}$ (the convergence of the full sequence follows from the uniqueness of the limit).

Proposition 3.6. Problem (1.6) admits a minimizer in the class $\mathcal{K}(\Omega)$.
Proof. Let $\left(K_{n}\right)_{n}$ be a minimizing sequence for Problem 1.6 in $\mathcal{K}(\Omega)$. Using Blaschke's theorem (see [3, Theorem 6.1]), we can find a compact connected proper subset $K$ of $\bar{\Omega}$ such that up to a subsequence, still denoted by the same index, $K_{n}$ converges to $K$ with respect to the Hausdorff distance. Then, of course $\partial \Omega \subset K$ and by Proposition 3.5, $u_{K_{n}}$ converges to $u_{K}$ strongly in $H_{0}^{2}(\Omega)$. Finally, thanks to the lower semicontinuity of $\mathcal{H}^{1}$ with respect to the Hausdorff convergence of connected sets, we deduce that $K$ is a minimizer of Problem (1.6).
3.3. Ahlfors regularity. We recall that a set $K \subset \mathbb{R}^{2}$ is said to be Ahlfors regular of dimension 1, if there exist some constants $c>0, r_{0}>0$ and $C>0$ such that for every $r \in\left(0, r_{0}\right)$ and for every $x \in K$ the following holds

$$
\begin{equation*}
c r \leq \mathcal{H}^{1}\left(K \cap B_{r}(x)\right) \leq C r . \tag{3.4}
\end{equation*}
$$

Note that for a closed connected nonempty set $K$ the lower bound in (3.4) is trivial: indeed, for all $x \in K$ and for all $r \in(0, \operatorname{diam}(K) / 2)$ we have $K \cap \partial B_{r}(x) \neq \emptyset$, and then

$$
\begin{equation*}
\mathcal{H}^{1}\left(K \cap B_{r}(x)\right) \geq r . \tag{3.5}
\end{equation*}
$$

Theorem 3.7. Let $\Omega \subset \mathbb{R}^{2}$ be a bounded Lipschitz domain. Let $(v, K)$ be a solution of Problem 3.3 with $\operatorname{diam}(K)>0$. Then $K$ is Ahlfors regular. More precisely, there exists $C>0$ and $r_{0}>0$ such that for all $x \in K$ and $0 \leq r \leq r_{0}$ we have

$$
\int_{B_{r}(x) \cap \Omega}|e(u)-G|^{2} d x+\mathcal{H}^{1}\left(K \cap B_{r}(x)\right) \leq 2 \pi r+C\|f\|_{p^{2}} r^{2} .
$$

Proof. It easily follows from the fact that $\Omega$ is a Lipschitz domain, that there exists $r_{0}>0$ such that $\partial B(x, r) \cap \bar{\Omega}$ is connected for all $x \in \Omega$ and $r \leq r_{0}$. Then to prove the Theorem, it suffice to compare ( $u, K$ ) with the admissible competitor ( $w, K_{r}$ ) defined by

$$
\begin{equation*}
K_{r}=\left(K \backslash B_{r}(x)\right) \cup\left(\partial B_{r}(x) \cap \bar{\Omega}\right), \tag{3.6}
\end{equation*}
$$

and $w=u \mathbf{1}_{\Omega \backslash B_{r}(x)}$. By this way we obtain

$$
\int_{B_{r}(x) \cap \Omega}|e(u)-G|^{2} d x+\mathcal{H}^{1}\left(K \cap B_{r}(x)\right) \leq \int_{B_{r}(x) \cap \Omega}|G|^{2} d x+2 \pi r,
$$

and finally using also Remark 3.4 we deduce that

$$
\int_{B_{r}(x) \cap \Omega}|e(u)-G|^{2} d x+\mathcal{H}^{1}\left(K \cap B_{r}(x)\right) \leq C\|f\|_{p} r^{2}+2 \pi r,
$$

which ends the proof.
3.4. $C^{1}$ regularity. The $C^{1}$ regularity of minimizers will follow from the following observation.

Proposition 3.8. If $(u, K)$ is a minimizer of Problem (3.3), then it is an almost-minimizer of the Grifith Energy. In other words there exists $C>0$ and $r_{0}>0$ such that for every competitor $\left(v, K^{\prime}\right)$ in the ball $B_{r}(x)$ and for all $r \leq r_{0}$ we have

$$
\int_{B_{r}}|e(u)|^{2} d x+\mathcal{H}^{1}\left(K \cap B_{r}(x)\right) \leq \int_{B_{r}}|e(v)|^{2} d x+\mathcal{H}^{1}\left(K^{\prime} \cap B_{r}(x)\right)+C r^{1+\frac{1}{2}}
$$

Proof. Let $\left(v, K^{\prime}\right)$ be a competitor for $(u, K)$ in the ball $B_{r}(x)$. Then since $u$ is a minimizer for Problem (3.3) and ( $v, K^{\prime}$ ) coincides with $(u, K)$ outside $B_{r}$ we have

$$
\int_{B_{r}}|e(u)-G|^{2} d x+\mathcal{H}^{1}\left(K \cap B_{r}\right) \leq \int_{B_{r}}|e(v)-G|^{2} d x+\mathcal{H}^{1}\left(K^{\prime} \cap B_{r}\right)
$$

which implies
$\int_{B_{r}}|e(u)|^{2} d x+\mathcal{H}^{1}\left(K \cap B_{r}\right) \leq \int_{B_{r}}|e(v)|^{2}-2 \int_{B_{r}} e(v): G d x+2 \int_{B_{r}} e(u): G d x+\mathcal{H}^{1}\left(K^{\prime} \cap B_{r}\right)$.
Now by elliptic regularity (see Remark 3.4) we have

$$
\int_{B_{r}}|G|^{2} d x \leq C r^{2}
$$

thus after Cauchy-Schwarz we arrive at
$\int_{B_{r}}|e(u)|^{2} d x+\mathcal{H}^{1}\left(K \cap B_{r}\right) \leq \int_{B_{r}}|e(v)|^{2}+C r\left(\int_{B_{r}}|e(v)|^{2}\right)^{\frac{1}{2}}+C r\left(\int_{B_{r}}|e(u)|^{2}\right)^{\frac{1}{2}}+\mathcal{H}^{1}\left(K^{\prime} \cap B_{r}\right)$.
Now by the proof of Ahlfors-Regularity we know that

$$
\int_{B_{r}}|e(u)-G|^{2} d x+\mathcal{H}^{1}\left(K \cap B_{r}\right) \leq C_{A} r
$$

which implies in particular that

$$
\int_{B_{r}}|e(u)|^{2} d x \leq 2 \int_{B_{r}}|e(u)-G|^{2} d x+2 \int_{B_{r}}|G|^{2} d x \leq 2 C_{A} r+C r^{2} \leq C^{\prime} r,
$$

provided that $r_{0} \leq 1$ (that we can assume). Returning to the inequality obtained before, we infer that

$$
\int_{B_{r}}|e(u)|^{2} d x+\mathcal{H}^{1}\left(K \cap B_{r}\right) \leq \int_{B_{r}}|e(v)|^{2}+C r\left(\int_{B_{r}}|e(v)|^{2}\right)^{\frac{1}{2}}+C r^{1+\frac{1}{2}}+\mathcal{H}^{1}\left(K^{\prime} \cap B_{r}\right) .
$$

We now divide the argument in two alternatives. Either

$$
\int_{B_{r}}|e(v)|^{2} \geq\left(C^{\prime}+C_{A}\right) r
$$

and then

$$
\int_{B_{r}}|e(u)|^{2} d x+\mathcal{H}^{1}\left(K \cap B_{r}\right) \leq\left(C^{\prime}+C_{A}\right) r \leq \int_{B_{r}}|e(v)|^{2}
$$

Or $\int_{B_{r}}|e(v)|^{2} \leq\left(C^{\prime}+C_{A}\right) r$ but then

$$
\int_{B_{r}}|e(u)|^{2} d x+\mathcal{H}^{1}\left(K \cap B_{r}\right) \leq \int_{B_{r}}|e(v)|^{2}+C r^{1+\frac{1}{2}}+\mathcal{H}^{1}\left(K^{\prime} \cap B_{r}\right) .
$$

In both cases we always have

$$
\int_{B_{r}}|e(u)|^{2} d x+\mathcal{H}^{1}\left(K \cap B_{r}\right) \leq \int_{B_{r}}|e(v)|^{2}+\mathcal{H}^{1}\left(K^{\prime} \cap B_{r}\right)+C r^{1+\frac{1}{2}} .
$$

which achieves the proof of the Proposition.
Therefore, by the results contained in [17] and [18], (see also [4]), we get the following theorem.
Theorem 3.9. Let $K \in \mathcal{K}(\Omega)$ be a solution for Problem (3.3). Then $K \cap \Omega$ is $C^{1, \alpha}$ outside a set of Hausdorff dimension strictly less than one.

## Appendix A. Quasi-EVERYWhere traces of $W^{k, 2}$ Functions

In the following statement we say that a property holds ( $m, p$ )-q.e. if it holds true outside a set of zero capacity of order $m \geq 1$ and integrability exponent $p$. We refer to [1] for the definition of capacity $\mathrm{Cap}_{m, p}$.

Let $K \subset \mathbb{R}^{N}$ be a closed set and $u \in W^{k, p}\left(\mathbb{R}^{N}\right)$. For $1<p<+\infty$ there exists a nice characterization of the space $W_{0}^{k, p}\left(\mathbb{R}^{N} \backslash K\right)$ in terms of traces of $u$ on $K$ as follows. If $\alpha$ is a multiindex such that $|\alpha| \leq k$ we say that $\left.\partial^{\alpha} u\right|_{K}=0$ if

$$
\lim _{r \rightarrow 0}\left(f_{B_{r}\left(x_{0}\right)}\left|\partial^{\alpha} u\right| d x\right)=0 \text { for }(k-|\alpha|, p) \text {-q.e. } x_{0} \text { on } K
$$

Then [1, Theorem 9.1.3.] says that

$$
\begin{equation*}
W_{0}^{k, p}\left(\mathbb{R}^{N} \backslash K\right)=\left\{u \in W^{k, p}\left(\mathbb{R}^{N}\right) \text { s.t. }\left.\partial^{\alpha} u\right|_{K}=0 \text { for all } 0 \leq \alpha \leq k-1\right\} \tag{A.1}
\end{equation*}
$$

where here $W_{0}^{k, p}\left(\mathbb{R}^{N} \backslash K\right)$ denotes the closure of $C_{c}^{\infty}\left(\mathbb{R}^{N} \backslash K\right)$ with respect to the $W^{k, p}\left(\mathbb{R}^{N}\right)$ norm. The following lemma is a generalization of the above fact, adapted to a domain $\Omega$ instead of $\mathbb{R}^{N}$. The proof uses an extension operator for $\Omega \subset \mathbb{R}^{N}$, i.e. a linear mapping $\mathcal{E}: W^{k, p}(\Omega) \rightarrow W^{k, p}\left(\mathbb{R}^{N}\right)$ such that $\mathcal{E}(u)=u$ in $\Omega$ and

$$
\|\mathcal{E}(u)\|_{W^{k, p}\left(\mathbb{R}^{2}\right)} \leq C\|u\|_{W^{k, p}(\Omega)}
$$

The construction of such operator for a Lipschitz domain $\Omega$, is classical (see for instance $[2$, Chapter 5]). By Lipschitz domain we mean a bounded open set whose boundary is locally the graph of a Lipschitz function at every point of the boundary.

The strategy for a domain $\Omega$ is then to apply the characterization in (A.1) to the extended function $\mathcal{E}(u)$ defined on $\mathbb{R}^{N}$, but it is not straightforward that the trace of $u$, which is a function defined only on $\Omega$, coincides with the trace of $\mathcal{E}(u)$ on the boundary $\partial \Omega$. For a Lipschitz domain this happens to be true thanks to a result in the book [16]. Here is then a general statement that we can prove with this strategy, that for simplicity we write only in the particular case $p=2$.
Lemma A.1. Let $\Omega \subset \mathbb{R}^{N}$ be a bounded Lipschitz domain and $K \subset \bar{\Omega}$ be a closed set. We consider the subspace of $W^{k, 2}(\Omega)$ defined by
$H_{0, K}^{k}(\Omega):=\left\{u \in W^{k, 2}(\Omega)\right.$ such that $\left.\partial^{\alpha} u\right|_{K}=0$ for all multiindices $\alpha$ such that $\left.0 \leq \alpha \leq k-1\right\}$, where by $\left.\partial^{\alpha} u\right|_{K}=0$ we mean that

$$
\lim _{r \rightarrow 0}\left(f_{B_{r}\left(x_{0}\right) \cap \Omega}\left|\partial^{\alpha} u\right| d x\right)=0 \text { for }(k-|\alpha|, 2) \text {-q.e. } x_{0} \text { on } K
$$

Then
(1) $H_{0, K}^{k}(\Omega)=\left\{u \in W^{k, 2}(\Omega)\right.$ such that $\left.\mathcal{E}(u) \in W_{0}^{k, 2}\left(\mathbb{R}^{N} \backslash K\right)\right\}$.
(2) $H_{0, K}^{k}(\Omega) \subset W^{k, 2}(\Omega)$ is closed for the strong topology of $W^{k, 2}(\Omega)$.
(3) $H_{0, K}^{k}(\Omega) \subset W^{k, 2}(\Omega)$ is closed for the weak topology of $W^{k, 2}(\Omega)$.
(4) If $u \in H_{0, K}^{k}(\Omega)$ then there exists a sequence $\varphi_{n} \in C_{c}^{\infty}\left(\mathbb{R}^{2}\right)$ such that $\operatorname{supp}\left(\varphi_{n}\right) \cap K=\emptyset$ for all $n$ and $\varphi_{n} \rightarrow u$ strongly in $W^{k, 2}(\Omega)$.

Proof. - Proof of (1). We begin with the proof of (1) that will actually be the key ingredient for all the other points. We already know by [1, Theorem 9.1.3.] that

$$
\begin{equation*}
W_{0}^{k, 2}\left(\mathbb{R}^{N} \backslash K\right)=\left\{u \in W^{k, 2}\left(\mathbb{R}^{N}\right) \text { s.t. }\left.\partial^{\alpha} u\right|_{K}=0 \text { for all } 0 \leq \alpha \leq k-1\right\} \tag{A.2}
\end{equation*}
$$

Thus to conclude, we only need to prove that for $(k-|\alpha|, 2)$-q.e. $x_{0} \in \partial \Omega$,

$$
\lim _{r \rightarrow 0}\left(f_{B_{r}\left(x_{0}\right) \cap \Omega}\left|\partial^{\alpha} u\right| d x\right)=\lim _{r \rightarrow 0}\left(f_{B_{r}\left(x_{0}\right)}\left|\partial^{\alpha} \mathcal{E}(u)\right| d x\right)
$$

This fact actually follows from [16, Proposition 2 page 207]. To say just a few words, the proof in [16] uses that $\partial^{\alpha}(\mathcal{E}(u))$, as a global Sobolev function on the full space $\mathbb{R}^{N}$, can be written as a convolution with a potential, namely $v:=\partial^{\alpha}(\mathcal{E}(u))=G_{\alpha} * f$ where $f \in L^{2}$. Then by the continuity behavior of Bessel potentials proved by Meyers in [20, Theorem 3.2.] we know that this type of function has a sort of pointwise continuity property, in the sense that

$$
v\left(x_{0}\right)=\lim _{\substack{x \rightarrow x_{0} \\ x \notin E}} v(x)
$$

where $E$ has a controlled capacity of the form $\operatorname{Cap}_{m, p}\left(E \cap B\left(x_{0}, r\right)\right)=o\left(r^{N-1}\right)$. From this we deduce that the limit of averages intersected with $\Omega$ or without intersection with $\Omega$ must coincide $(k-|\alpha|, 2)$-q.e. We refer to [16, Proposition 2 page 207] for more details. This achieves the proof of (1), and we can notice that it does not depend on the choice of the extension operator. Notice also that the argument used in Proposition A. 2 could give an independent proof for the special case $k=1$.

- Proof of (2). Let $u_{n}$ be a sequence in $H_{0, K}^{k}(\Omega)$ such that $u_{n} \rightarrow u$ in $W^{k, 2}(\Omega)$. Since the extension operator $\mathcal{E}$ is continuous on $W^{k, 2}$ it follows that $\mathcal{E}\left(u_{n}\right) \rightarrow \mathcal{E}(u)$ in $W^{k, 2}\left(\mathbb{R}^{2}\right)$ and by use of (1) we know that $\mathcal{E}\left(u_{n}\right) \in W_{0}^{k, 2}\left(\mathbb{R}^{2} \backslash K\right)$, for all $n$. But now by definition, $W_{0}^{k, 2}\left(\mathbb{R}^{N} \backslash K\right)$ is a closed subspace of $W^{k, 2}\left(\mathbb{R}^{N} \backslash K\right)$ and since $\mathcal{E}\left(u_{n}\right) \rightarrow \mathcal{E}(u)$ we obtain $\mathcal{E}(u) \in W_{0}^{k, 2}\left(\mathbb{R}^{2} \backslash K\right)$. By applying (1) again we deduce that $u \in H_{0, K}^{k}(\Omega)$.
- Proof of (3). Let $u_{n} \in H_{0, K}^{k}(\Omega)$ be a sequence that converges weakly in $W^{k, 2}(\Omega)$ to some limit $u$, in other words $\left\langle u_{n}, \varphi\right\rangle \rightarrow\langle u, \varphi\rangle$ for all $\varphi \in W^{k, 2}(\Omega)$, where the brackets means the scalar product in $W^{k, 2}(\Omega)$. Since $H_{0, K}^{k}(\Omega)$ is a closed subset of $W^{k, 2}(\Omega)$ (for the strong topology), it follows that $H_{0, K}^{k}(\Omega)$ is itself a Hilbert space endowed with the same norm and scalar product of $W^{k, 2}(\Omega)$. We can therefore extract a subsequence that converges for the weak topology in $H_{0, K}^{k}(\Omega)$ to some limit function $h$ that must still belong to $H_{0, K}^{k}(\Omega)$. This says in particular that $\left\langle u_{n}, \varphi\right\rangle \rightarrow\langle h, \varphi\rangle$ for all $\varphi \in H_{0, K}^{k}(\Omega)$. By uniqueness of the limit in the weak topology of $H_{0, K}^{k}(\Omega)$, we must have $h=u$, concluding that finally $u \in H_{0, K}^{k}(\Omega)$ as desired.
- Proof of (4). This item is a direct consequence of (1).

As seen in the proof of Lemma A.1, if $\Omega$ is a Lipschitz domain and $\mathcal{E}$ is an extension operator for $W^{1,2}(\Omega)$, then it holds true that for every Sobolev function $u \in W^{1,2}(\Omega)$,

$$
\begin{equation*}
\lim _{r \rightarrow 0}\left(f_{B_{r}\left(x_{0}\right) \cap \Omega}|u| d x\right)=\lim _{r \rightarrow 0}\left(f_{B_{r}\left(x_{0}\right)}|u| d x\right) \quad \text { for }(1,2)-\text { q.e. } x_{0} \in \partial \Omega \tag{A.3}
\end{equation*}
$$

In the paper we would need a more uniform result of the same kind, but where $u$ is not fixed but could also depend on $r$ and converges strongly in $W^{1,2}$. Since we could not find in the literature a short proof for this property, we write in this appendix a complete argument. For that purpose we use an explicit extension operator, i.e. a linear mapping $\mathcal{E}: W^{1,2}(\Omega) \rightarrow W^{1,2}\left(\mathbb{R}^{2}\right)$ such that
$\mathcal{E}(u)=u$ in $\Omega$ and

$$
\begin{equation*}
\|\mathcal{E}(u)\|_{W^{1,2}\left(\mathbb{R}^{2}\right)} \leq C\|u\|_{W^{1,2}(\Omega)} . \tag{A.4}
\end{equation*}
$$

There are several ways of constructing such an operator. Here we follow the approach in [13] as follows. If $\partial \Omega$ locally coincides around $x_{0} \in \partial \Omega$ with the graph of the Lipschitz mapping $\gamma: \mathbb{R}^{N-1} \rightarrow \mathbb{R}$, and if $u$ is compactly supported in a cylinder of size $h>0$ centered at $x_{0}$, one can use the direct formula

$$
\mathcal{E}(u)\left(x^{\prime}, x_{N}\right)=\left\{\begin{array}{cl}
u\left(x^{\prime}, x_{N}\right) & \text { if } x_{N}>\gamma\left(x^{\prime}\right)  \tag{A.5}\\
u\left(x^{\prime}, 2 \gamma\left(x^{\prime}\right)-x_{N}\right) & \text { if } x_{N}<\gamma\left(x^{\prime}\right) .
\end{array}\right.
$$

For the general case one can cover the boundary $\partial \Omega$ with a finite union of cylinders of same size $h$ and use a partition of unity, leading to an extension operator of the form

$$
\begin{equation*}
\mathcal{E}(u)=\sum_{i} \theta_{i} \mathcal{E}_{i}(u), \tag{A.6}
\end{equation*}
$$

where $\mathcal{E}_{i}$ is of the form (A.5) in a suitable local coordinate system. We refer to Theorem 1 in section 4.4. of [13] for further detail, where in particular the estimate (A.4) is established.

Now we focus on the following proposition, and its immediate corollary which is used only in the proof of Lemma 2.11.

Proposition A.2. Let $\Omega$ be a bounded Lipschitz domain and $\mathcal{E}$ be the extension operator defined in (A.6). Let $r \rightarrow 0$ and $\left(u_{r}\right)$ be a sequence in $W^{1,2}(\Omega)$ such that $u_{r} \rightarrow u \in W^{1,2}(\Omega)$. Let $x_{0} \in \partial \Omega$. Then $\lim _{r \rightarrow 0} f_{B(x, r)} \mathcal{E}\left(u_{r}\right)(y) d y=0$ implies $\lim _{r \rightarrow 0} f_{B(x, r) \cap \Omega} u_{r}(y) d y=0$.

Remark A.3. As an immediate corollary, and under the same assumptions of Proposition A.2, we have

$$
\lim _{r \rightarrow 0} f_{B(x, r)}\left|\mathcal{E}\left(u_{r}\right)(y)\right| d y=0 \Longrightarrow \lim _{r \rightarrow 0} f_{B(x, r) \cap \Omega}\left|u_{r}(y)\right| d y=0
$$

Compare with (A.3).
Proof. We first assume that $\partial \Omega \cap B\left(x_{0}, r_{0}\right)$ coincides with a portion of the graph of a Lipschitz function $\gamma: \mathbb{R}^{N} \rightarrow \mathbb{R}$, that $\Omega \cap B\left(x_{0}, r\right)=B\left(x_{0}, r\right) \cap\left\{\left(x^{\prime}, x_{N}\right): x_{N}>\gamma\left(x^{\prime}\right)\right\}$ for all $r<r_{0}$ and that $\mathcal{E}(u)=u\left(x^{\prime}, 2 \gamma\left(x^{\prime}\right)-x_{N}\right)$ in $\left\{x_{N}<\gamma\left(x^{\prime}\right)\right\}$. Then we define, for $r<r_{0}$,

$$
B_{r}^{+}=B\left(x_{0}, r\right) \cap\left\{x_{N}>\gamma\left(x^{\prime}\right)\right\} \quad \text { and } \quad B_{r}^{-}=B\left(x_{0}, r\right) \cap\left\{x_{N}<\gamma\left(x^{\prime}\right)\right\},
$$

and write

$$
\int_{B_{r}\left(x_{0}\right)} \mathcal{E}(u) d x=\int_{B_{r}^{+}} u d x+\int_{B_{r}^{-}} u\left(x^{\prime}, 2 \gamma\left(x^{\prime}\right)-x_{N}\right) d x^{\prime} d x_{N} .
$$

On the other hand by a simple change of variable wich is linear in the $x_{N}$ variable, together with Fubini's Theorem, we see that

$$
\int_{B_{r}^{-}}\left|u\left(x^{\prime}, 2 \gamma\left(x^{\prime}\right)-x_{N}\right)\right| d x^{\prime} d x_{N}=\int_{A_{r}}|u(x)| d x
$$

where $A_{r}$ is the "reflected" domain $\Phi\left(B_{r}^{-}\right)$, with $\Phi\left(x^{\prime}, x_{N}\right):=\left(x^{\prime}, 2 \gamma\left(x^{\prime}\right)-x_{N}\right)$. Since $\gamma$ is Lipschitz with a constant depending only on $\Omega$, we infer that there exists $\Lambda, \lambda>0$ such that $B_{\lambda r}^{+} \subset A_{r} \subset B_{\Lambda r}^{+}$for all $r \leq r_{0}$. We can therefore estimate the difference between the average in
$A_{r}$ and the one in $B_{R}^{+}$in the following way, denoting by $R=\Lambda r$, (the constant $C$ below depends only on $\Omega$ and can change from line to line)

$$
\begin{align*}
\left|f_{A_{r}} u d x-f_{B_{R}^{+}} u d x\right| & =\left|f_{A_{r}}\left(u-f_{B_{R}^{+}} u\right) d x\right| \leq f_{A_{r}}\left|u-f_{B_{R}^{+}} u\right| d x \\
& \leq C r^{-2} \int_{B_{R}^{+}}\left|u-f_{B_{R}^{+}} u\right| d x \leq C r^{-1}\left(\int_{B_{R}^{+}}\left(u-f_{B_{R}^{+}} u\right)^{2} d x\right)^{\frac{1}{2}} \\
& \leq C\left(\int_{B_{R}^{+}}|\nabla u|^{2} d x\right)^{\frac{1}{2}} . \tag{A.7}
\end{align*}
$$

For the last inequality we have used the fact that the domains $B_{R}^{+}$are uniformly Lipschitz as $r$ is going to 0 , in order to apply a Poincaré inequality with uniform constant in all of the $B_{R}^{+}$. Now we apply the above to $u_{r} \in W^{1,2}(\Omega)$, instead of $u$ and we know that

$$
\int_{B_{R}^{+}}\left|\nabla u_{r}\right|^{2} d x \xrightarrow{R \rightarrow 0} 0,
$$

because $u_{r} \rightarrow u$ strongly in $W^{1,2}(\Omega)$, and this proves that

$$
\begin{equation*}
\left|f_{A_{r}} u_{r} d x-f_{B_{R}^{+}} u_{r} d x\right| \xrightarrow{r \rightarrow 0} 0 . \tag{A.8}
\end{equation*}
$$

Now assume that $f_{B\left(x_{0}, r\right)} \mathcal{E}\left(u_{r}\right) d x \xrightarrow{r \rightarrow 0} 0$. Notice that $\left|A_{r}\right|=\left|B_{r}^{-}\right|$because the jacobian of $\Phi$ is equal to 1 . Then

$$
\begin{aligned}
2 f_{B_{r}^{+}} u_{r} d x & =\frac{1}{\left|B_{r}^{+}\right|}\left(\int_{B_{r}^{+}} u_{r} d x+\int_{A_{r}} u_{r} d x+\int_{B_{r}^{+}} u_{r} d x-\int_{A_{r}} u_{r} d x\right) \\
& =\frac{\left|B\left(x_{0}, r\right)\right|}{\left|B_{r}^{+}\right|} f_{B\left(x_{0}, r\right)} \mathcal{E}\left(u_{r}\right) d x+\frac{\left|B_{r}^{-}\right|}{\left|B_{r}^{+}\right|}\left(f_{B_{r}^{+}} u_{r} d x-f_{A_{r}} u_{r} d x\right)+\left(1-\frac{\left|B_{r}^{-}\right|}{\left|B_{r}^{+}\right|}\right) f_{B_{r}^{+}} u_{r} d x
\end{aligned}
$$

So finally
$\left(1+\frac{\left|B_{r}^{-}\right|}{\left|B_{r}^{+}\right|}\right) f_{B_{r}^{+}} u_{r} d x=\frac{\left|B\left(x_{0}, r\right)\right|}{\left|B_{r}^{+}\right|} f_{B\left(x_{0}, r\right)} \mathcal{E}\left(u_{r}\right) d x+\frac{\left|B_{r}^{-}\right|}{\left|B_{r}^{+}\right|}\left(f_{B_{r}^{+}} u_{r} d x-f_{A_{r}} u_{r} d x\right) \xrightarrow{r \rightarrow 0} 0$
because of (A.8). This proves that $f_{B\left(x_{0}, r\right) \cap \Omega} u_{r} d x=f_{B_{r}^{+}} u_{r} d x \rightarrow 0$ and finishes the proof of the proposition in the particular case when $\mathcal{E}(u)$ coincides with the formula $u\left(x^{\prime}, 2 \gamma\left(x^{\prime}\right)-x_{N}\right)$ under the graph of the Lipschtiz function $\gamma$.

In the general case $\partial \Omega$ is covered by a finite number of Lipschitz graphs $\gamma_{i}$ and $\mathcal{E}(u)$ is of the form $\sum_{i} \theta_{i} \mathcal{E}_{i}(u)$ where $\theta_{i} \in C_{c}^{\infty}\left(\mathbb{R}^{2}\right)$ is a partition of unity and $\mathcal{E}_{i}$ is the extension of $u$ relatively to the graph $\gamma_{i}$. Applying the above argument to each of the $\mathcal{E}_{i}(u)$ we conclude as follows.

We assume that

$$
\lim _{r \rightarrow 0}\left(f_{B_{r}\left(x_{0}\right)} \mathcal{E}\left(u_{r}\right) d x\right)=0
$$

which in other words says

$$
\lim _{r \rightarrow 0}\left(f_{B_{r}\left(x_{0}\right)} \sum_{i} \theta_{i} \mathcal{E}_{i}\left(u_{r}\right) d x\right)=0
$$

or again after change of variable,

$$
\lim _{r \rightarrow 0} f_{B_{r}^{+}} u_{r} d x+\sum_{i}\left(f_{A_{r}^{i}} \theta_{i} \circ \Phi_{i}^{-1} u_{r} d x\right)=0
$$

where $A_{r}^{i}$ is the reflexion of $B_{r}^{-}$relatively to the graph of $\gamma_{i}$, i.e. $A_{r}^{i}=\Phi_{i}\left(B_{r}^{-}\right)$with $\Phi_{i}\left(x^{\prime}, x_{N}\right):=$ $\left(x^{\prime}, 2 \gamma_{i}\left(x^{\prime}\right)-x_{N}\right)$ (here the coordinate system $\left(x^{\prime}, x_{N}\right)$ should be taken relatively to the one associated with $\gamma_{i}$ ). Now we denote by $I$ the set of indices $i$ for which $\theta_{i}$ is not identically zero around $x_{0}$, and we let $N$ be the cardinal of $I$ (which should actually be at most two).

Now by construction $\Phi_{i}\left(x_{0}\right)=x_{0}$ for all $i$ and $\Phi_{i}$ is Lipschitz with constant depending only on $\Omega$. Since $\theta_{i}$ is a smooth function we deduce that

$$
\begin{equation*}
\left|\theta_{i} \circ \Phi_{i}^{-1}(x)-\theta_{i}\left(x_{0}\right)\right| \leq C\left|x-x_{0}\right| \tag{A.9}
\end{equation*}
$$

Let $\Lambda>0$ be such that $A_{r}^{i} \subset B_{\Lambda r}^{+}$for all small $r$. Then since $u_{r} \in L^{2}(\Omega)$ we can use (A.9) to obtain

$$
\begin{equation*}
\left(f_{A_{r}^{i}}\left|\theta_{i} \circ \Phi_{i}^{-1}(x)-\theta_{i}\left(x_{0}\right)\right|\left|u_{r}\right| d x\right) \leq C\left(\int_{B\left(x_{0}, \Lambda r\right)}\left|u_{r}\right|^{2} d x\right)^{\frac{1}{2}} \xrightarrow{r \rightarrow 0} 0 \tag{A.10}
\end{equation*}
$$

because $u_{r}$ converges in $L^{2}$ to $u$. Next, following an argument similar to the one of (A.7), denoting by $R:=\Lambda r$, using (A.10) and the fact that $\sum_{i} \theta_{i}\left(x_{0}\right)=1$ we obtain some function $e(r) \rightarrow 0$ for which

$$
\begin{aligned}
\left|f_{B_{r}^{-}} \mathcal{E}\left(u_{r}\right) d x-f_{B_{R}^{+}} u_{r} d x\right|= & \left|\left(\sum_{i \in I} f_{A_{r}^{i}} \theta_{i} \circ \Phi_{i}^{-1} u_{r} d x\right)-f_{B_{R}^{+}} u_{r} d x\right| \\
\leq & \left|\sum_{i} \theta_{i}\left(x_{0}\right) f_{A_{r}^{i}}\left(u_{r}-f_{B_{R}^{+}} u_{r}\right) d x\right|+e(r) \\
\leq & C r^{-2} \int_{B_{R}^{+}}\left|u_{r}-f_{B_{R}^{+}} u_{r}\right| d x+e(r) \\
\leq & C\left(\int_{B_{R}^{+}}\left|\nabla u_{r}\right|^{2} d x\right)^{\frac{1}{2}}+e(r) \\
& \xrightarrow{r \rightarrow 0} 0 .
\end{aligned}
$$

Then as before we can write

$$
\left(1+\frac{\left|B_{r}^{-}\right|}{\left|B_{r}^{+}\right|}\right) f_{B_{r}^{+}} u_{r} d x=\frac{\left|B\left(x_{0}, r\right)\right|}{\left|B_{r}^{+}\right|} f_{B\left(x_{0}, r\right)} \mathcal{E}\left(u_{r}\right) d x+\frac{\left|B_{r}^{-}\right|}{\left|B_{r}^{+}\right|}\left(f_{B_{r}^{+}} u_{r} d x-f_{B_{r}^{-}} \mathcal{E}\left(u_{r}\right) d x\right) \xrightarrow{r \rightarrow 0} 0
$$

which finishes the proof in the general case, and thus the proposition is now proved.
We end this section with variants of the Poincaré inequality related to the traces of functions on positive capacity subsets. The first result can be compared to [1, Corollary 8.2.2].

Corollary A.4. Let $\Omega \subset \mathbb{R}^{2}$ be a bounded Lipschitz domain and let the closed set $K \subset \bar{\Omega}$ be such that $\operatorname{Cap}_{1,2}(K)>0$. Then there exists a constant $C>0$ such that for all $u \in H_{0, K}^{1}(\Omega)$,

$$
\|u\|_{L^{2}(\Omega)} \leq C\|\nabla u\|_{L^{2}(\Omega)}
$$

where $H_{0, K}^{k}(\Omega)$ is the space defined in Lemma A.1.

Proof. We prove the result by contradiction. Assume that there exists a sequence $u_{k} \in H_{0, K}^{1}(\Omega)$ such that

$$
\left\|u_{k}\right\|_{L^{2}(\Omega)} \geq k\left\|\nabla u_{k}\right\|_{L^{2}(\Omega)} .
$$

By renormalizing the sequence, we can assume that $\left\|\nabla u_{k}\right\|_{L^{2}(\Omega)} \rightarrow 0$, as $k \rightarrow \infty$, while $\left\|u_{k}\right\|_{L^{2}(\Omega)}=1$. It follows that, passing to a subsequence, $u_{k}$ converges strongly to a constant $a_{0} \neq 0$. But it immediately follows from Lemma A. 1 that $a_{0} \equiv 0$ on $K$ quasi-everywhere, which contradicts $a_{0} \neq 0$ as a constant.

The second result concerns uniform constants for the Poincaé inequality on thin domains $\Omega^{h}$, where the trace of the functions vanish on subsets of the form $K_{h} \times\{0\} \subset \Omega \times\{0\}$. It should be compared with [19, Theorem D.1].

Theorem A.5. Let $\Omega$ be a bounded Lipschitz domain, $\Omega^{h}:=\Omega \times(0,1)$, and let

$$
\mathbf{A}_{K}^{h}:=\left\{u \in W^{1,2}\left(\Omega^{h}\right) ;\left.u\right|_{K_{h} \times\{0\}}=0\right\} .
$$

Then there exists $h_{0}>0$ such that for a constant $C>0$ uniform in $h<h_{0}$ we have

$$
\forall u \in \mathbf{A}_{K}^{h} \quad\|u\|_{L^{2}\left(\Omega^{h}\right)} \leq C\|\nabla u\|_{L^{2}\left(\Omega^{h}\right)} .
$$

We will provide a sketch of the proof, based on [19, Theorem D.1], adapted to our situation. We will follow the steps leading to Lemma 2.6, replacing the rigidity estimate in Corollary 2.3 by the Poincaré inequality applied to the Lipschitz domains:

$$
\left\|u-a_{x^{\prime}}\right\|_{L^{2}\left(\mathcal{Q}_{j}\left(x^{\prime}, h\right)\right)} \leq C h^{2}\|\nabla u\|_{L^{2}\left(\mathcal{Q}_{j}\left(x^{\prime}, h\right)\right)}
$$

with $a_{x^{\prime}}=0$ whenever $K_{c h} \cap \overline{\Phi_{j}^{-1}\left(S_{\xi^{\prime}, h}\right)} \neq \emptyset$. We can then conclude with the following lemma parallel to Lemma 2.7:

Lemma A.6. Let $\Omega, \Omega^{h}, K$ be as defined above. Then there exist constants $h_{0}>0, C>0$, $0<\bar{c}<1$, depending only on $\Omega$ and $K$, such that given $h<h_{0}, u \in \mathbf{A}_{K}^{h}$, there exists a scalar function $a: \Omega \rightarrow \mathbb{R}$ which vanishes on $K_{\bar{c} h}$ such that the estimates

$$
\begin{equation*}
\|u-a\|_{L^{2}\left(\Omega^{h}\right)}^{2} \leq C h^{2}\|\nabla u\|_{L^{2}\left(\Omega^{h}\right)}^{2}, \quad\|\nabla a\|_{L^{2}(\Omega)}^{2} \leq \frac{C}{h}\|\nabla u\|_{L^{2}\left(\Omega^{h}\right)}^{2}, \tag{A.12}
\end{equation*}
$$

and

$$
\begin{equation*}
\|a\|_{L^{\infty}(\Omega)}^{2} \leq \frac{C}{h}\|\nabla u\|_{L^{2}\left(\Omega^{h}\right)}^{2}, \tag{A.13}
\end{equation*}
$$

hold true.
Proof of Theorem A.5. Let $u \in \mathbf{A}_{K}^{h}$ and let $a \in W^{1,2}(\Omega)$ be chosen according to Lemma A.6. Note that $a \in H_{0, K}^{1}(\Omega)$, and hence by Corollary A. 4 we obtain for a uniform constant $C$

$$
\|a\|_{L^{2}(\Omega)}^{2} \leq C\|\nabla a\|_{L^{2}(\Omega)}^{2},
$$

which yields, through the second inequality in (A.12)

$$
\|a\|_{L^{2}\left(\Omega^{h}\right)}^{2}=h\|a\|_{L^{2}(\Omega)}^{2} \leq C h\|\nabla a\|_{L^{2}(\Omega)}^{2} \leq C\|\nabla u\|_{L^{2}\left(\Omega^{h}\right)}^{2} .
$$

The conclusion follows from the first inequality in (A.12).

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(M. R. Pakzad) Laboratoire Imath, Université de Toulon, CS 60584-83041, TOULON CEDEX 9, France

Email address: pakzad@univ-tln.fr
(A. Lemenant) Institut universitaire de France (IUF) et Université de Lorraine, CNRS, IECL, F-54000 Nancy, France

Email address: antoine.lemenant@univ-lorraine.fr


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